

A catalogue of Structural And Morphological Measurements for DES Y1

F. Tarsitano^{1*}, W. G. Hartley²,

A. Amara¹, A. Bluck³, C. Bruderer¹, M. Carollo³, C. Conselice⁴, P. Melchior⁵, B. Moraes², A. Refregier¹, I. Sevilla-Noarbe⁶, J. Woo³, T. M. C. Abbott⁷, S. Allam⁸, J. Annis⁸, S. Avila⁹, M. Banerji^{10,11}, E. Bertin^{12,13}, D. Brooks², D. L. Burke^{14,15}, A. Carnero Rosell^{16,17}, J. Carretero¹⁸, C. E. Cunha¹⁴, C. B. D'Andrea¹⁹, L. N. da Costa^{16,17}, C. Davis¹⁴, J. De Vicente⁶, S. Desai²⁰, P. Doel², J. Estrada⁸, J. Frieman^{8,21}, J. García-Bellido²², D. Gruen^{14,15}, R. A. Gruendl^{23,24}, G. Gutierrez⁸, D. Hollowood²⁵, K. Honscheid^{26,27}, D. J. James²⁸, T. Jeltema²⁵, E. Krause^{29,30}, K. Kuehn³¹, N. Kuropatkin⁸, O. Lahav², M. A. G. Maia^{16,17}, F. Menanteau^{23,24}, R. Miquel^{32,18}, A. A. Plazas³⁰, A. K. Romer³³, A. Roodman^{14,15}, E. Sanchez⁶, B. Santiago^{34,16}, R. Schindler¹⁵, M. Smith³⁵, R. C. Smith⁷, M. Soares-Santos³⁶, F. Sobreira^{37,16}, E. Suchyta³⁸, M. E. C. Swanson²⁴, G. Tarle³⁹, D. Thomas⁹, V. Vikram⁴⁰, A. R. Walker⁷

(DES Collaboration)

¹ Institute for Particle Physics and Astrophysics, ETH Zurich, Wolfgang-Pauli-Strasse 27, CH-8093 Zurich, Switzerland

² Department of Physics & Astronomy, University College London, Gower Street, London, WC1E 6BT, UK

³ Department of Physics, ETH Zurich, Wolfgang-Pauli-Strasse 16, CH-8093 Zurich, Switzerland

⁴ University of Nottingham, School of Physics and Astronomy, Nottingham NG7 2RD, UK

⁵ Department of Astrophysical Sciences, Princeton University, Peyton Hall, Princeton, NJ 08544, USA

⁶ Centro de Investigaciones Energéticas, Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain

⁷ Cerro Tololo Inter-American Observatory, National Optical Astronomy Observatory, Casilla 603, La Serena, Chile

⁸ Fermi National Accelerator Laboratory, P. O. Box 500, Batavia, IL 60510, USA

⁹ Institute of Cosmology & Gravitation, University of Portsmouth, Portsmouth, PO1 3FX, UK

¹⁰ Institute of Astronomy, University of Cambridge, Madingley Road, Cambridge CB3 0HA, UK

¹¹ Kavli Institute for Cosmology, University of Cambridge, Madingley Road, Cambridge CB3 0HA, UK

¹² CNRS, UMR 7095, Institut d'Astrophysique de Paris, F-75014, Paris, France

¹³ Sorbonne Universités, UPMC Univ Paris 06, UMR 7095, Institut d'Astrophysique de Paris, F-75014, Paris, France

¹⁴ Kavli Institute for Particle Astrophysics & Cosmology, P. O. Box 2450, Stanford University, Stanford, CA 94305, USA

¹⁵ SLAC National Accelerator Laboratory, Menlo Park, CA 94025, USA

¹⁶ Laboratório Interinstitucional de e-Astronomia - LIneA, Rua Gal. José Cristino 77, Rio de Janeiro, RJ - 20921-400, Brazil

¹⁷ Observatório Nacional, Rua Gal. José Cristino 77, Rio de Janeiro, RJ - 20921-400, Brazil

¹⁸ Institut de Física d'Altes Energies (IFAE), The Barcelona Institute of Science and Technology, Campus UAB, 08193 Bellaterra (Barcelona) Spain

¹⁹ Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, PA 19104, USA

²⁰ Department of Physics, IIT Hyderabad, Kandi, Telangana 502285, India

²¹ Kavli Institute for Cosmological Physics, University of Chicago, Chicago, IL 60637, USA

²² Instituto de Física Teórica UAM/CSIC, Universidad Autonoma de Madrid, 28049 Madrid, Spain

²³ Department of Astronomy, University of Illinois at Urbana-Champaign, 1002 W. Green Street, Urbana, IL 61801, USA

²⁴ National Center for Supercomputing Applications, 1205 West Clark St., Urbana, IL 61801, USA

²⁵ Santa Cruz Institute for Particle Physics, Santa Cruz, CA 95064, USA

²⁶ Center for Cosmology and Astro-Particle Physics, The Ohio State University, Columbus, OH 43210, USA

²⁷ Department of Physics, The Ohio State University, Columbus, OH 43210, USA

²⁸ Harvard-Smithsonian Center for Astrophysics, Cambridge, MA 02138, USA

²⁹ Department of Astronomy/Steward Observatory, 933 North Cherry Avenue, Tucson, AZ 85721-0065, USA

³⁰ Jet Propulsion Laboratory, California Institute of Technology, 4800 Oak Grove Dr., Pasadena, CA 91109, USA

³¹ Australian Astronomical Observatory, North Ryde, NSW 2113, Australia

³² Institució Catalana de Recerca i Estudis Avançats, E-08010 Barcelona, Spain

³³ Department of Physics and Astronomy, Pevensey Building, University of Sussex, Brighton, BN1 9QH, UK

³⁴ Instituto de Física, UFRGS, Caixa Postal 15051, Porto Alegre, RS - 91501-970, Brazil

³⁵ School of Physics and Astronomy, University of Southampton, Southampton, SO17 1BJ, UK

³⁶ Brandeis University, Physics Department, 415 South Street, Waltham MA 02453

³⁷ Instituto de Física Gleb Wataghin, Universidade Estadual de Campinas, 13083-859, Campinas, SP, Brazil

³⁸ Computer Science and Mathematics Division, Oak Ridge National Laboratory, Oak Ridge, TN 37831

³⁹ Department of Physics, University of Michigan, Ann Arbor, MI 48109, USA

⁴⁰ Argonne National Laboratory, 9700 South Cass Avenue, Lemont, IL 60439, USA

ABSTRACT

We present a structural and morphological catalogue for 45 million objects selected from the first year data of the Dark Energy Survey (DES). Single Sérsic fits and non-parametric measurements are produced for g , r and i filters. The parameters from the best-fitting Sérsic model (total magnitude, half-light radius, Sérsic index, axis ratio and position angle) are measured with GALFIT; the non-parametric coefficients (concentration, asymmetry, clumpiness, Gini, M20) are provided using the Zurich Estimator of Structural Types (ZEST+). To study the statistical uncertainties, we consider a sample of state-of-the-art image simulations with a realistic distribution in the input parameter space and then process and analyse them as we do with real data: this enables us to quantify the observational biases due to PSF blurring and magnitude effects and correct the measurements as a function of magnitude, galaxy size, Sérsic index (concentration for the analysis of the non-parametric measurements) and ellipticity. We present the largest structural catalogue to date: we find that accurate and complete measurements for all the structural parameters are typically obtained for galaxies with SExtractor MAG_AUTO_I ≤ 21 . Indeed, the parameters in the filters i and r can be overall well recovered up to MAG_AUTO ≤ 21.5 , corresponding to a fitting completeness of $\sim 90\%$ below this threshold, for a total of 25 million galaxies. The combination of parametric and non-parametric structural measurements makes this catalogue an important instrument to explore and understand how galaxies form and evolve. The catalogue described in this paper will be publicly released alongside the Dark Energy Survey collaboration Y1 cosmology data products.

Key words: galaxy evolution, galaxy morphology, galaxy structure

1 INTRODUCTION

Any explanation of the formation and evolution of galaxies must necessarily include a description of the diverse forms that galaxies take. The morphology of the luminous components of a galaxy, including its classification or decomposition into a bulge and disk (e.g., Kormendy 1977; de Jong 1996) or identification of features such as bars, rings or lenses (e.g., Kormendy 1979; Combes & Sanders 1981; Elmegreen et al. 1996), are a result of its aggregated formation history. Assigning meaningful morphological types or quantifying the distribution of light across the extent of a population of galaxies, is therefore of fundamental importance in understanding the processes that govern their evolution.

A quantitative description of galaxy morphology is typically expressed in terms of structural parameters (brightness, size, shape) and properties of the light distribution (concentration, asymmetry and clumpiness), though human classifications are still used. The development of *citizen science* projects like Galaxy Zoo (Lintott et al. 2008; Simmons et al. 2017; Willett et al. 2017) and sophisticated machine learning algorithms (Lahav et al. 1995; Lahav 1995; Huertas-Company et al. 2008, 2015; Banerji et al. 2010; Dieleman et al. 2015) have helped to maintain the relevance of these perception-based morphologies in the current literature. Nevertheless, most recent work on the subject of galaxy morphologies rely on either *parametric* or *non-parametric* approaches to quantify the galaxy’s light distribution.

Parametric methods consist of a two-dimensional fitting of the flux intensity of the galaxy, including parametric mathematical models of the light fall-off and deconvolution of the point spread function (PSF) from the observed galaxy image. The most general assumed function for this purpose is the Sérsic profile (Sérsic 1963). The second class performs an analysis of the light distribution within a certain elliptical area, usually defined through the Petrosian radius associated to the galaxy. Common estimates are of the degree to

which the light is concentrated, quantifying the asymmetry of the light distribution and searching for clumpy regions: this method is called *CAS system* (Concentration, Asymmetry and Smoothness or Clumpiness) and can be extended with further parameters, Gini and M20 (Conselice 2003; Abraham et al. 2003; Lotz et al. 2004; Law et al. 2007). These parameters together can describe the major features of galaxy structure without resorting to model assumptions about the galaxy’s underlying form, as is done with the Sérsic profile. However, they are determined without a PSF deconvolution and need an additional calibration.

Even alone, distributions of morphological quantities represent powerful constraints on possible galaxy formation scenarios. But combined with other physical quantities, they can provide key insights into the processes at play, supporting or even opening new ideas on evolutionary mechanisms (Kauffmann et al. 2004; Weinmann et al. 2006; Schawinski et al. 2007; van der Wel 2008a,b; Bamford et al. 2009; Schawinski et al. 2014). For instance, the relationship between the masses, luminosities and sizes of massive disks and spheroids suggests dissipative formation processes within hierarchical dark matter assembly (White & Rees 1978; Fall & Efstathiou 1980) or the occurrence of galaxy-galaxy mergers (Toomre & Toomre 1972; Toomre 1977; Barnes 1988; Naab & Burkert 2003; Conselice 2003; Lin et al. 2004; Conselice 2008; Conselice et al. 2008; Jogee 2009; Jogee et al. 2009). Similarly, bulges, disks and bars may be formed by secular evolution processes (Kormendy 1979; Kormendy & Kennicutt 2004; Bournaud et al. 2007; Genzel et al. 2008; Fisher & Drory 2008; Sellwood 2014) or by the interplay between smooth and clumpy cold streams and disk instabilities (Dekel et al. 2009a,b). In this sense bulges may be formed without major galaxy mergers, as is often thought.

Of particular interest in recent years, have been the questions over the degree to which galaxy environment impacts upon morphology (e.g. Dressler 1980; Postman et al. 2005; Lani et al. 2013; Kuutma et al. 2017), and the connection between morphology and cessation of star formation in galaxies (e.g. Blanton et al. 2003; Martig et al. 2009; Bell et al. 2012; Woo et al. 2015). Faced with often

* federica.tarsitano@phys.ethz.ch

subtle correlations or hidden variables within strong correlations, these questions demand far greater statistical power and measurement precision than had been possible from the available data sets in the preceding decades. These demands require efficient pipelines to automate and streamline the analysis of large astronomical data sets. GALAPAGOS (Gray et al. 2009; Häußler et al. 2011; Barden et al. 2012) is perhaps the most widely used of such pipelines. It offers a routine to simplify the process of source detection, to cut postage stamps, prepare masks for neighbours if needed and estimate a robust sky background and has been used at both low redshift in the GEMS survey (Häussler et al. 2007), and at higher redshift on the CANDELS (van der Wel et al. 2012) data.

At low redshift the state-of-the-art to date are the catalogues constructed from Sloan Digital Sky Survey (SDSS, York et al. 2000) data, in particular the bulge+disk catalogue of Simard et al. (2011) numbering almost 1 million galaxies. Such statistical power has been lacking at higher redshifts, but the advent of large-scale cosmology experiments optimised for weak lensing analyses, such as the Dark Energy Survey (DES) and Hyper Suprime-Cam (HSC) (Miyazaki et al. 2012), provide a great opportunity to fill in much of this gap. DES is the largest galaxy survey to date, with a narrower PSF and images typically two magnitudes deeper than the SDSS.

In order to create as complete a set of structural measurements for DES as possible we adopt both parametric and non-parametric approaches, using the software GALFIT (Peng et al. 2002, 2010) for parametric Sérsic fitting and ZEST+ for a non-parametric analysis of the structural properties of our galaxy sample. The first provides us with the measurements of the magnitude, effective radius, Sérsic Index, axis ratio and orientation angle of the galaxy; the second one outputs an extended set of parameters, completing the CAS system with Gini and M20, plus the values of magnitude, half light radius and ellipticity, measured within the galaxy Petrosian ellipse.

The scale of the DES data set requires a new dedicated pipeline in order to handle the DES data structure, optimise run-time performance, automate the process of identifying and handling neighbouring sources and prepare tailored postage stamps for input to the two software packages. The resulting dataset is by far the largest catalogue of structural parameters measured to date, numbering 45 million galaxies, which exceeds previous catalogues by more than an order of magnitude in size, and reaches redshift, $z \sim 1$. It includes parametric and non-parametric measurements in three photometric bands, intended to be used in concert and to provide a comprehensive view of the galaxies' morphologies. In this sense, our DES Y1 catalogue constitutes a significant step in our capabilities to study the nature of galaxy morphology in the Universe.

This paper is structured as follows: in Section 2 we give an overview of the Dark Energy Survey, describing the data and the image simulation data we used for this work. In Section 3 we focus on the details of our sample selection and pre-fitting routine, presenting the algorithms developed to prepare and process the data. Sections 4 and 5 are dedicated to the parametric and non-parametric fits, respectively. In each of these two sections, we present a detailed description of the fitting software used for this work, discuss the completeness and validation of the fitted sample from each method, provide an overview of the characteristics of the catalogue and perform a calibration of the output quantities with image simulations. The calibration for the i band are shown in those sections; Appendix A includes the calibration maps also for the g and r filters. Section 5 also introduces a set of basic cuts as a starting point in building a science-ready sample. Finally in Section 7 we summarise our work. A manual explaining the catalogue columns is presented in Appendix B.

2 DATA

2.1 The Dark Energy Survey

The Dark Energy Survey (DES) (DES Collaboration 2005; The DES Collaboration 2016) is a wide-field optical imaging survey covering 5000 deg² of the southern equatorial hemisphere in $grizY$ bands¹. Survey observations began in August 2013 and over five years it will provide images of 300 million of galaxies up to redshift ~ 1.4 (Diehl et al. 2014). The survey is designed to have a combination of area, depth and image quality optimized for cosmology, and in particular the measurement of weak gravitational lensing shear. However, its rich data set is well-suited to many areas of astronomy, including galaxy evolution, Milky Way and Local Group science, stellar populations and Solar System science (Abbott et al. 2016). DES uses the Dark Energy Camera (DECam), a mosaic imager with a 2.2° diameter field of view and a pixel scale of 0.263'' per pixel mounted at the prime focus of the Victor M. Blanco 4m Telescope at Cerro Tololo Inter-American Observatory. During the requested 525 observing nights it is expected to reach photometric limits of $g = 24.6$, $r = 24.4$, $i = 23.7$, $z = 22.7$ and $Y = 21.5$ (10σ limits in 1.5'' apertures assuming 0.9'' seeing) following ten single-epoch exposures of 90 seconds each for $griz$ and 45 seconds each for Y (Flaugher 2005).

The DES data are processed, calibrated and archived through the DES Data Management (DESDM) system (Drlica-Wagner et al. 2017; Morganson et al. 2018), consisting of an image processing pipeline which performs image de-trending, astrometric calibration, photometric calibration, image co-addition and SExtractor (Bertin & Arnouts 1996) catalogue creation. The DESDM imaging co-addition combines overlapping single-epoch images in a given filter, which are then remapped to artificial tiles in the sky so that one co-add image per band is produced for every tile. These tiles are padded to ensure that each object is entirely contained in at least one tile, but also results in a small fraction of duplicate objects found in different tiles which are removed at a later stage. In order to account for PSF variations caused by object location in the focal plane and the combination of images with different seeing, the catalogue creation process uses PSFex (Bertin 2011, 2013) to model the PSF. PSFex produces a basis set of model components that are combined via linear combination into a location-dependent PSF. The final step combines the photometry of each co-add object into a single entry in multi-wavelength SExtractor catalogues. For more details about the DESDM co-addition and PSF modelling we refer the reader to Sevilla et al. 2011, Desai et al. 2012 and Mohr et al. 2012.

In this work we use the DES Y1A1 COADD OBJECTS data release, comprising 139,142,161 unique objects spread over about 1800 deg² in 3707 co-add tiles, constructed from the first year of DES survey operations. The tiles are combinations of 1-5 exposures in each of the $grizY$ filters and the average coverage depth at each point in the retained footprint is ~ 3.5 exposures. We consider 3690 tiles in total: the catalogue for the remaining tiles, located in the 30 deg² of cadenced supernovae fields, will be presented in future work. The data include all the products of the DESDM pipeline and imaging co-addition (the co-add tiles and their respective segmentation maps, the PSF models and the SExtractor catalogues), plus the Y1A1 GOLD catalogue (Drlica-Wagner et al. 2017). In the Y1A1 GOLD catalogue, the data collected in DES year-one have been char-

¹ <http://www.darkenergysurvey.org>

4 DES Collaboration

acterised and calibrated in order to form a sample which minimises the occurrence of artefacts and systematic features in the images. It further provides value-added quantities such as the star-galaxy classifier MODEST and photo- z estimates. GOLD magnitudes are corrected for interstellar extinction using stellar locus regression (SLR) (High et al. 2009). We combine the SExtractor DESDM catalogues with the Y1A1 GOLD catalogue to make the sample selection, as described in section 3.1, and we also benefit from the application of the MODEST classifier during the analysis of the completeness of our fitting results, reported in more detail in section 4.2.

2.2 Image simulation data

In fitting galaxy light profiles, faint magnitude regimes are well known to present larger systematic errors in the recovered galaxy sizes, fluxes and ellipticities (Bernstein et al. 2002; Häussler et al. 2007; Melchior & Viola 2012). A larger FWHM of the PSF can also introduce increased uncertainties and systematic errors during morphological estimation. In order to overcome these issues we use sophisticated image simulations to derive multi-parameter vectors that quantify any biases arising from our analyses, data quality or modelling assumptions. The simulations we use for this purpose are produced by the *Ultra Fast Image Generator* (UFIG) (Bergé et al. 2013) run on the Blind Cosmology Challenge simulation (BCC, Busha et al. 2013) and released for DES Y1 as UFIG-BCC. UFIG-BCC covers an area of 1750 deg^2 and includes images which are calibrated to match the DES Y1 instrumental effects, galaxy distribution and survey characteristics. Briefly, an input catalogue of galaxies is generated based on the results of an N-body simulation with an algorithm to reproduce the observed luminosity and colour-density relations.

3 PRE-FITTING PIPELINE

In this section we describe first the sample selection we apply to the DES Y1A1 COADD OBJECTS, discussing the cuts applied and the initial distributions. Then we describe the process which prepares the data to be fitted both with parametric and non-parametric approach.

3.1 Sample Selection

For this work we use a tile-by-tile approach, independently for each filter: every step from the sample selection itself to the fitting process is performed separately in each tile and band, with the exception of an overall i -band magnitude cut and fiducial star-galaxy separation. We organise the Y1A1 GOLD catalogue into sub-catalogues to include the objects in each co-add tile and match them with the relevant DESDM SExtractor catalogues, extracted from that tile. We apply cuts to specific flags in the catalogues and to the parameters we use as priors for the fits in order to remove the most probable point-like sources, whilst avoiding removing galaxies. In addition we remove a small amount of the survey area in order to work with objects whose SExtractor detection and images are reliable and well-suited for the fitting process. An object is selected if it fulfils the following requirements:

- $\text{FLAGS}_X = 0$;
- $\text{GOLD_MAG_AUTO}_I \leq 23$;
- $\text{FLUX_RADIUS}_X > 0$;
- $\text{KRON_RADIUS}_X > 0$;

SELECTION TYPE	SELECTION CUT
Gold match	$\text{IN_GOLD} = \text{True}$
Image flags	$\text{FLAGS}_x = 0$
S/G	$\text{CLASS_STAR}_i \leq 0.9$
Magnitude	$\text{MAG_AUTO}_i \leq 23$
Size (I)	$\text{FLUX_RADIUS} > 0 \text{ px}$
Size (II)	$\text{KRON_RADIUS} > 0 \text{ px}$
Regions	$\text{FLAGS_BADREGION} = 0$

Table 1. Summary of the cuts applied to the overlapping sample between the catalogue provided by the DESDM pipeline and the Y1A1 GOLD catalogue. The selected objects must satisfy the requirements described in section 3.1. x identifies the filter ($x = g, r, i$).

- $\text{CLASS_STAR}_I < 0.9$;
- $\text{FLAGS_BADREGION} = 0$,

where $X = g, r, i, z, Y$. The cut in FLAGS removes objects that are either saturated, truncated or have been de-blended. We apply the cuts using the i band as our reference band; indeed the seeing FWHM in this filter is on average the smallest of the five bands. In using the CLASS_STAR classifier at this stage we perform a conservative star-galaxy discrimination (S/G), so that we attempt a fit for any object which could be a galaxy. During the validation analysis we will remove further objects, applying a stricter classifier, named MODEST. We refer to section 4.2.1 for its definition and more details about its impact on this work. By GOLD_MAG_AUTO we refer to the SExtractor quantity MAG_AUTO, corrected by photometric calibration through SLR as provided by the Y1A1 GOLD catalogue (Drllica-Wagner et al. 2017). In the following sections we will simply use the original uncalibrated SExtractor MAG_AUTO. FLAGS_BADREGION is a flag from the Y1A1 GOLD catalogue tracing the objects that lie in problematic areas, which are close to high-density stellar regions and/or present ghosts and glints. The sample selection cuts described above are summarised in Table 1; the distributions of the variables considered during these initial cuts, comparing the selected sample with the entire dataset (in grey), are shown in Fig. 1.

3.2 Data processing

The co-add data used in this work are processed in a dedicated pre-fitting pipeline, called Selection And Neighbours Detection (SAND), which has been developed in order to prepare the postage stamps to be fit, their ancillary files in the formats required by GALFIT and ZEST+ and perform essential book-keeping operations. The pipeline performs three steps: *sample selection* (as described in section 3.1), *stamp cutting* and *identification of neighbouring sources*. It is important to note that the objects excluded by our initial sample selection (section 3.1) are still fit as neighbouring objects where appropriate. For this reason dedicated flags are assigned to each object in the sample, in order to trace their CLASS_STAR classification and possible anomalies in their photometric and structural properties. Collectively, we refer to these flags as STATUS_FLAGS, and document the components and possible values in Appendix B.

For each selected object, an image postage stamp is created,

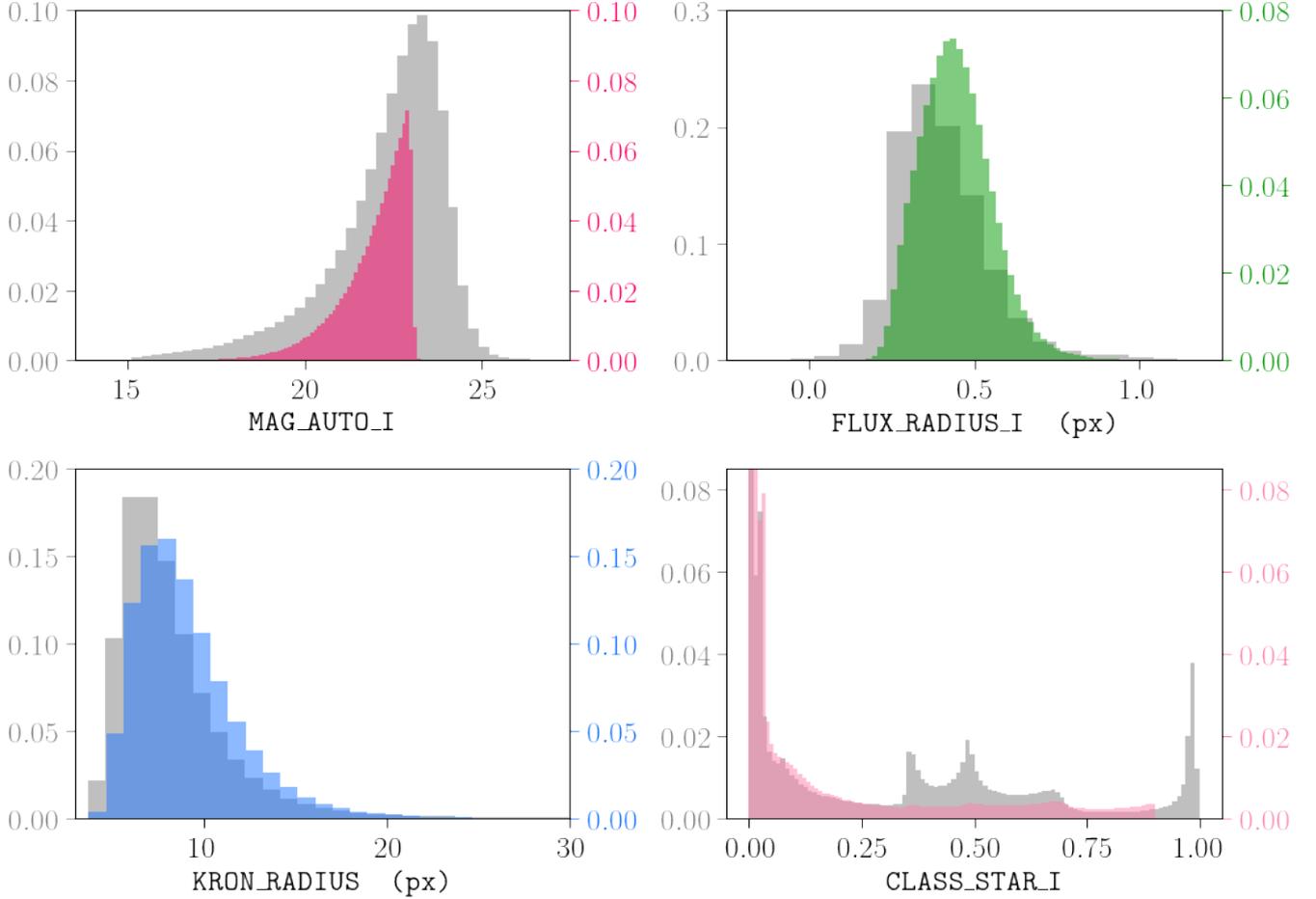


Figure 1. Distributions of the variables involved in the sample selection in the *i* band. From upper left to bottom right: MAG_AUTO, CLASS_STAR, FLUX_RADIUS and KRON_RADIUS. The cuts applied to each variable are described in more detail in section 3.1 and summarised in Table 1. In each panel the grey histogram refers to the whole dataset, while the coloured one represents the distribution in that variable for the selected sample.

282 initially with half-width equal to 3 times its Kron radius². Using the
283 relevant segmentation map, the algorithm calculates the percentage
284 of pixels that are not associated with sources (i.e. are background
285 pixels) and approves the stamp if the sky fraction is at least 60%.
286 Otherwise, the image stamp is rejected and is enlarged in size in
287 integer multiples of Kron radius until this requirement is satisfied.

288 The last step of the pre-fitting routine is dedicated to the identifi-
289 cation and cataloguing of neighbours: using the postage seg-
290 mentation maps it locates the neighbouring objects and, with the
291 above mentioned STATUS_FLAGS, identifies nearby potential stars
292 and/or galaxies with unreliable SExtractor detection. With this
293 last expression we refer to the objects which have unphysical SEx-
294 tractor parameters (negative sizes, magnitude set to standard error
295 values) and/or are flagged as truncated or saturated objects. In ad-
296 dition to their coordinates and SExtractor properties, the routine
297 catalogues the relative SExtractor magnitude and the presence
298 of overlapping Kron-like isophotes between the central galaxy and
299 its neighbours: these cases are then classified with two dedicated
300 flags, called ELLIPSE_FLAGS and MAX_OVERLAP_PERC, which are

301 fully described in Appendix B³. This information is now easily
302 accessible during the parametric fitting routine and helps to make
303 decisions on the models to be used to simultaneously fit the objects
304 lying in each stamp (see section 4.1); indeed, they are crucial also
305 to the non-parametric approach, since they communicate to ZEST+
306 all the necessary information to clean the neighbours in the stamps
307 and prepare them for the measuring routine which is described in
308 section 5.1.

309 4 PARAMETRIC FITS

310 4.1 GALFIT Setup

Image cutouts and PSF models appropriate to each individual object
are provided to GALFIT, which is used to find the best-fitting Sérsic
models. As reported in (Peng et al. 2002, 2010), the adopted Sérsic
function has the following form:

$$\Sigma(r) = \Sigma_e \exp \left\{ \left[-k \left(\frac{r}{R_e} \right)^{\frac{1}{n}} - 1 \right] \right\}, \quad (1)$$

³ By *Kron-like isophote* we refer to the Kron ellipse enlarged by a factor of 1.5.

² i.e. SExtractor KRON_RADIUS \times A_IMAGE.

where Σ_e is the pixel magnitude at the half-light radius R_e . The Sérsic index n quantifies the profile concentration: if n is large, we have a steep inner profile with a highly extended outer wing; inversely, when n is small, the inner profile is shallow and presents a steep truncation at large radii. In the case of $n = 1$ we have an exponential light profile. GALFIT produces measurements for the free parameters of the Sérsic function: central position, integrated magnitude (m_{tot}), effective radius (R_e) measured along the major axis, Sérsic index (n), axis ratio (q) and position angle (PA). The integrated magnitude is determined through its definition as a function of the flux (F_{tot}) integrated out to $r = \infty$ for the Sérsic profile:

$$m_{tot} = -2.5 \log \left(\frac{F_{tot}}{t_{exp}} \right) + mag_zpt, \quad (2)$$

where t_{exp} is the exposure time and mag_zpt is the zero-point magnitude, both indicated in the image header.

Apart from the central position, which is allowed to vary by only ± 1 pixel by a GALFIT constraints file, all the parameters are left free without constraints: for those, initial guesses are taken from the SExtractor DESDM catalogues (the exception being Sérsic index, which is always started at $n = 2$). Thanks to the large background area available in each stamp (pre-validated with the SAND algorithm), GALFIT is left free to estimate the background level⁴. The information provided by the SAND routine is adopted in order to optimise the simultaneous fitting procedure of the central galaxy and its neighbours. Using the ELLIPSE_FLAGS (introduced in section 3.2) it is easy to identify most of the neighbours, including faint companions, nearby stars, close objects with overlapping isophotes and neighbours with unreliable priors due to unphysical SExtractor measurements.

Companion objects three magnitudes fainter than the main galaxy are not fit. In the presence of overlapping isophotes, the relevant neighbouring object is fit simultaneously with the target galaxy (even in the cases where it is centred outside the stamp). However, if the overlapping region is 50% or larger than the area within the Kron-like ellipse occupied by the central galaxy, then although a fit is attempted, it is not considered for the analysis discussed in this paper. Given k_1 and k_2 as the effective Kron Radii of the central galaxy and its neighbour respectively, they are used to define the isophotes of those objects, intended as enlarged Kron-like ellipses. If the isophotes are not overlapping, but separated by less than the maximum between k_1 and k_2 , then the neighbour is fit simultaneously. Otherwise it is masked. If the neighbour is a star ($CLASS_STAR \geq 0.9$), it is simultaneously fit with a PSF model. Finally, if the stamp contains one or more neighbours whose initial guesses from SExtractor contain errors (for example negative magnitudes and radii), no fit is attempted. We adopt a Single Sérsic model with all its parameters free for neighbours also.

4.2 Fitting Completeness

GALFIT uses a non-linear least-squares algorithm which iterates χ^2 minimization in order to find the best solution given a large parameter space. However even when the algorithm outputs a solution, there could be cases where the estimation of one or more parameters is affected by numerical convergence issues, which makes the solution itself an unreliable and non-unique result. These cases include

⁴ During initial tests on the fitting routine we randomly selected a sub-sample of objects to be fitted with the background fixed to zero. The outcome of this test was that this choice does not change significantly the results.

correlated parameters, local minima and mathematically degenerate solutions (Peng et al. 2010, Section 6). GALFIT labels the affected parameters enclosing them in between stars (**). In such cases we classify the fit as non-converged and do not trust the set of structural parameters it provides.

We determine the fraction of converged and non-converged fits and investigate their properties and location in the DES field. We present our analysis for all filters taking the i band as reference to discuss the fitting properties and possible causes of failure and incompleteness.

We evaluate the fitting completeness by calculating the percentage of converged fits in differential bins of 0.2 magnitude. The completeness (C) is calculated by normalising the number of converged fits in each magnitude bin ($N(c|mag)$) to the number of objects which passed the sample selection (described in section 3.1) in that bin, as expressed in the following definition:

$$C_{|mag} = \frac{N(c|mag)}{N(c|mag) + N(nc|mag) + N(f|mag)}, \quad (3)$$

where $N(nc|mag)$ and $N(f|mag)$ refer to the fractions of non-converged and failed fits in each magnitude bin, respectively. We also derive the percentage of converged fits calculated within limiting magnitudes.

The results of this analysis are shown in Fig. 2. In the upper left inset (Panel A) the solid lines represent the fitting completeness in magnitude bins and the dashed lines the magnitude limited completeness. They are colour-coded by filter: green and orange lines refer to g and r band, respectively; brown and black to the i band. We start our discussion from the latter.

The dashed black line shows the completeness determined for a sample with a conservative star-galaxy (S/G) cut ($CLASS_STAR > 0.9$): the trend shows that $\sim 90\%$ of the fits are successful at magnitude ~ 17 , after which this value starts to decline and reaches $\sim 80\%$ at magnitude ~ 21 . The completeness decreases more rapidly towards fainter magnitudes.

The brown line shows the completeness after applying a star-galaxy cut based on the SPREAD_MODEL parameter. In this way a completeness of $\sim 85\%$ is reached at magnitude 21.5. More details about the star-galaxy classifier and the analysis are described in the next paragraph. We match the information given by in the first panel with the map in Panel B: each point represents a DES tile and is colour-coded by the percentage of converged fits in that tile. The grey region, with $100 < ra < 60$ and $-70 < dec < -58$, has been excluded from the sample selection because in the GOLD catalogue it is flagged due to its vicinity to the Large Magellanic Cloud (LMC). We observe that the regions with a higher percentage of non-converged fits are located at the East and West borders of the footprint, towards the Galactic plane. These regions are characterized by high stellar density, as shown in Pieres et al. 2017. One possibility is that many of the unconverged fits at relatively bright magnitudes are stellar contaminants and so there is a poorer completeness where the stellar spatial frequency is higher. Another scenario could be that the edges of the footprints were observed under poorer conditions, for instance with poorer seeing.

We now investigate the correlations between our fitting completeness and maps of survey characteristics (as introduced in Drlica-Wagner et al. 2017), and discuss the likely causes of failures, encompassing stellar contamination, the effect of PSF width, poor signal-to-noise and the effects of neighbouring sources.

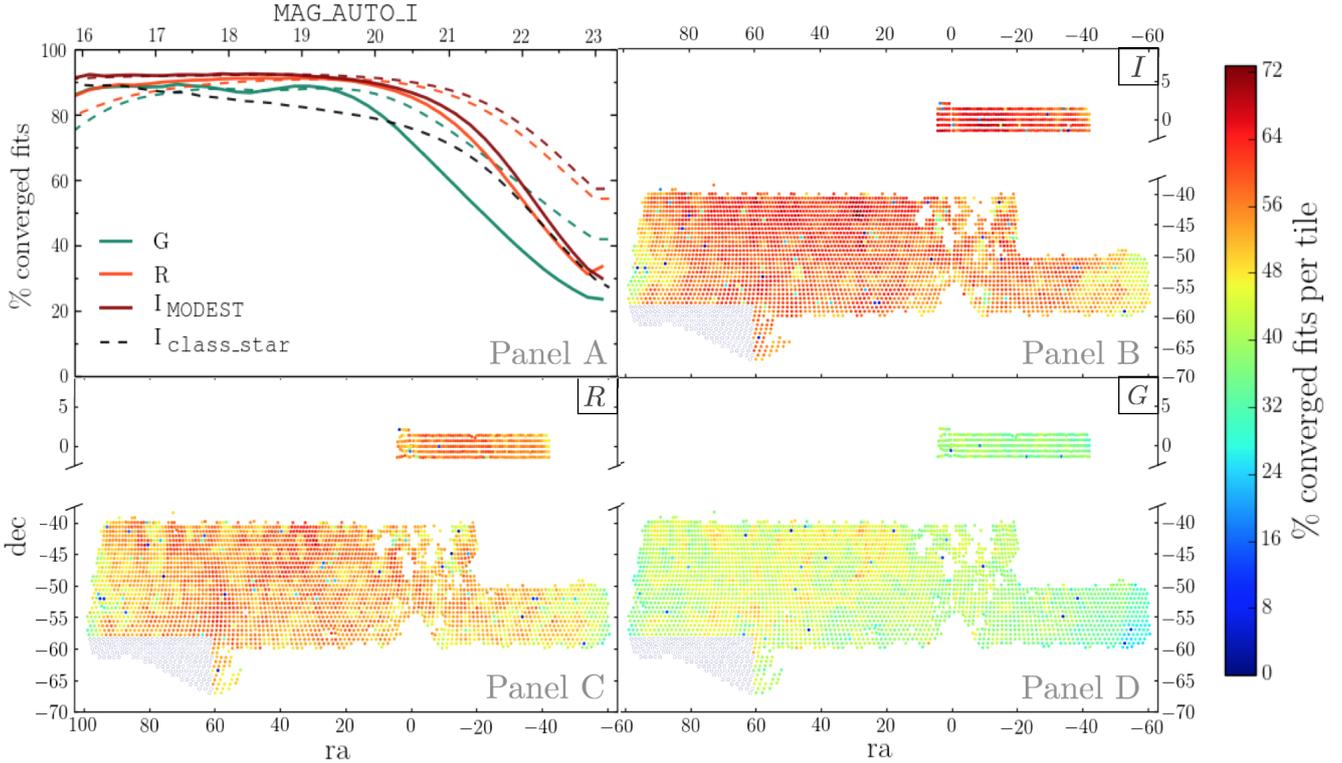


Figure 2. Panel A: fitting completeness in g , r and i bands (green, orange and brown lines, respectively), following star-galaxy separation using the MODEST classifier (see Section 4.2.1). The completeness, defined in eq. 3, is expressed in terms of the percentage of converged fits calculated in bins of 0.2 magnitude. Solid lines show the completeness in differential magnitude bins, while the dashed lines show results for magnitude-limited samples. The dashed black line shows the trend for the i band when using only a conservative S/G cut ($\text{CLASS_STAR} > 0.9$). Using the MODEST classifier we find that the completeness is 90% up to magnitude 21. Panels B, C, D: maps of the percentage of converged fits in g , r and i band in each tile (at $\text{mag_auto_}i < 23$). The grey region is entirely flagged as unsuitable for extra-galactic work due to its vicinity to the Large Magellanic Cloud (LMC). The regions with lower density of converged fits are found towards the Galactic Plane and close to the LMC. In the g band the percentage of converged fits is poorer, as expected, due to an overall broader PSF.

387 4.2.1 Stellar contamination

388 We used the neural network star-galaxy (S/G) classifier, included
389 as part of SExtractor, for a conservative initial criterion of star-
390 galaxy separation. We apply the cut $\text{CLASS_STAR} < 0.9$, in order
391 to remove only the most obvious stars, and to allow a user to per-
392 form their own S/G separation. Point sources will most likely fail to
393 achieve a converged solution in GALFIT and we therefore expect that
394 a substantial fraction of the incompleteness at bright magnitudes
395 seen in the black dotted line in Fig. 2 (panel A) is due to contami-
396 nation by stellar sources. This expectation is supported by the fact
397 that the regions with the lowest percentage of converged fits (Fig. 2,
398 panels B-D) are located in regions of known high stellar density.
399 Further, in the upper panel of Fig. 3 it can be seen that the converged
400 fraction at $i < 21.5$ depends strongly on the stellar density for the
401 CLASS_STAR S/G separation.

In Drlica-Wagner et al. (2017) it is shown that a simple cut in the SExtractor parameters SPREAD_MODEL and SPREADERR_MODEL can achieve a galaxy completeness of $\geq 98\%$, with $\leq 3\%$ stellar contamination at $i < 22$. This cut is known as MODEST classifier. SPREAD_MODEL is a morphological quantity which compares the source to both the local PSF and a PSF-convolved exponential model (Desai et al. 2012; Soumagnac et al. 2015). In order to optimise the separation of point-like and spatially extended sources, we use the i band as the reference band for object classification due to the

depth and superior PSF in this filter. The separation is defined via a linear combination of the SPREAD_MODEL and its uncertainty, the SPREADERR_MODEL:

$$\text{SPREAD_MODEL} + n \times \text{SPREADERR_MODEL} > thr, \quad (4)$$

402 where the coefficients $n = 1.67$ and $thr = 0.005$ are chosen as the
403 optimal compromise between the completeness and purity of the
404 galaxy sample. With the MODEST classifier we recover more than
405 $\sim 90\%$ converged fits at magnitude 20 and $\sim 85\%$ at magnitude
406 21.5.

407 We apply this additional S/G classification henceforth, and show
408 the converged fraction of galaxies under this additional classifica-
409 tion by the coloured lines in Fig. 2 and the black points in Fig. 3.
410 The dependence of converged fraction on stellar density is vastly
411 reduced with the SPREAD_MODEL classifier (though still present)
412 with a threefold increase in stellar density, from 0.5 to 1.5 stars per
413 sq. arcmin, causing just a 7% point drop in converged fraction. This
414 decrease is almost entirely explained by the expected contamination
415 rate of 3%.

416 4.2.2 PSF width

Obtaining deconvolved light profiles of galaxies with observed sizes close to the size of the PSF requires very accurate knowledge of the PSF. Errors in the PSF model can easily result in attempted fits

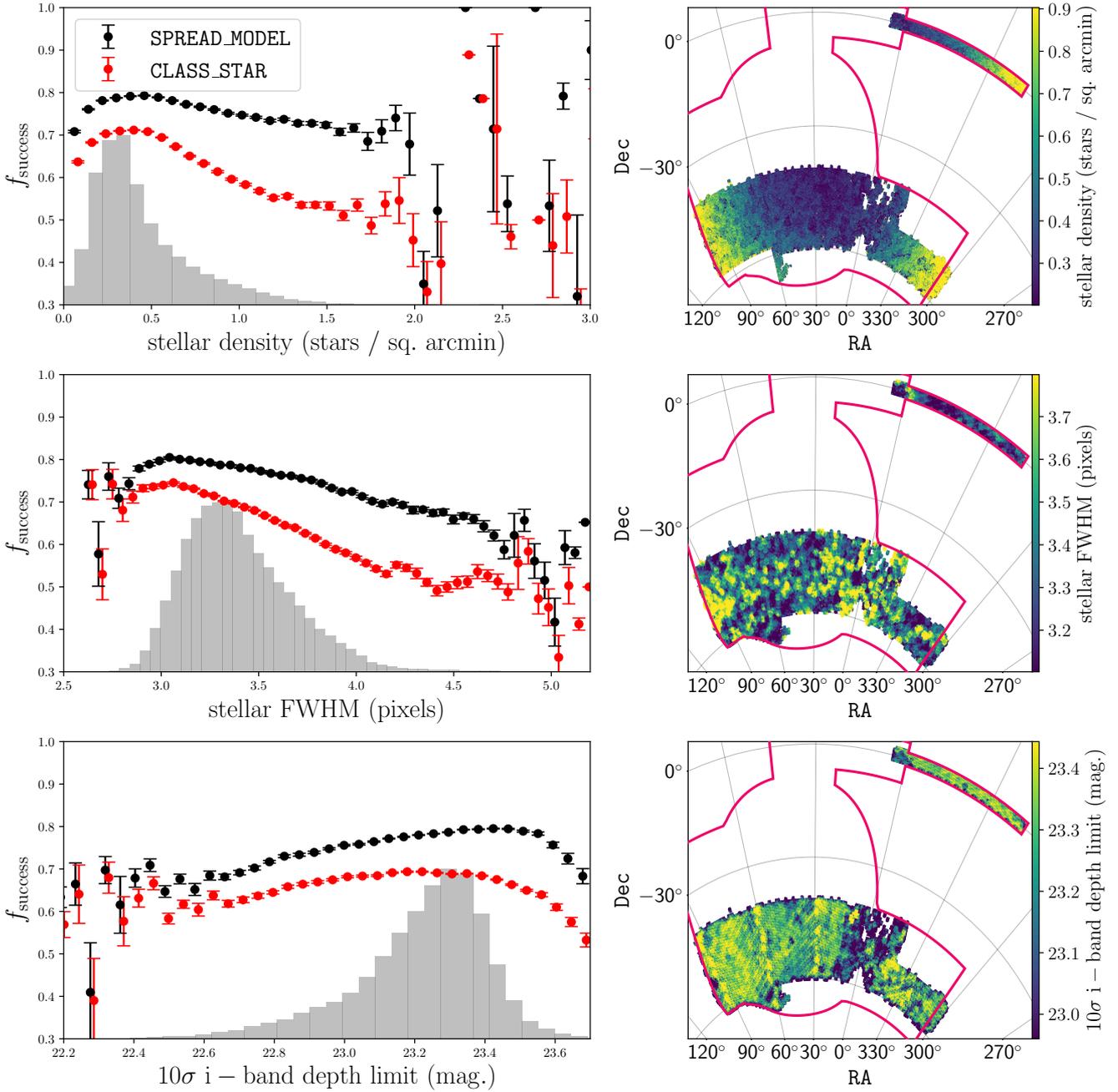


Figure 3. Dependence of fitting completeness at $i < 21.5$ on spatially-dependent survey characteristics, stellar density, PSF FWHM and i -band image depth (top, middle and bottom panels respectively). Grey histograms show the relative distributions of the characteristics in terms of survey area. The results for the galaxy sample are shown, following two star-galaxy classifiers: SExtractor CLASS_STAR (red points) and an additional criterion based on SPREAD_MODEL (black points, see text). Uncertainties are derived by bootstrap resampling. After the improved S/G separation, the fitting completeness is only weakly dependent on survey characteristics, and a high completeness ($> 80\%$) can be maintained with only minimal loss of area. The results at $i < 22$ are very similar in terms of the correlations with survey characteristics, but with overall lower converged fraction.

not converging, or in biased parameters (see section 4.4). Here, we assess the fitting incompleteness due to the varying PSF width across the DES survey area. We calculate the completeness for different sub-populations of the sample, identified with certain values of the ratio between the galaxy size estimated by SExtractor and the PSF

size; we indicate this parameter with ξ , defined as follows:

$$\xi = \frac{\text{FLUX_RADIUS}}{\text{PSF_radius}}, \quad (5)$$

where we calculate the size of the PSF as the radius of the circular aperture enclosing half of the flux of the PSF itself. The left panel in Fig. 4 shows the completeness calculated in bins of 1 magnitude for five different populations: $\xi \leq 0.75$, $0.75 < \xi \leq 1$, $1 < \xi \leq 1.25$,

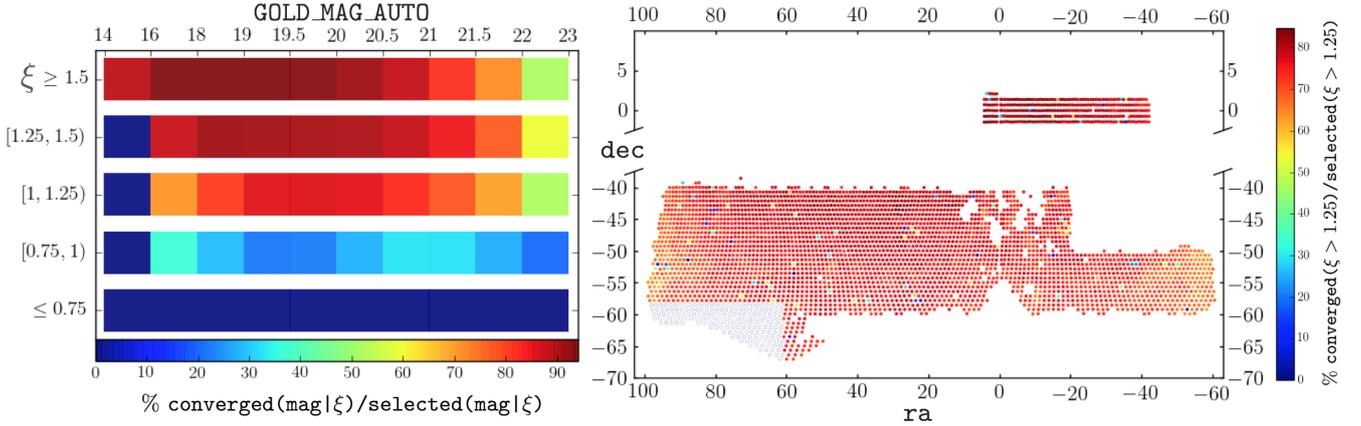


Figure 4. **Left panel:** fitting completeness calculated in differential bins of magnitude. The sample is divided into sub-populations, according to different ranges of the parameter $\xi = \text{FLUX_RADIUS}/\text{PSF_radius}$, as reported on the y-axis. Each population is represented by a bar, colour-coded by the percentage of converged fits in each magnitude bin. The figure shows that failed fits are more frequent for the objects with size smaller than the PSF or comparable with it. A critical drop occurs for the population with $\xi < 1.25$. **Right panel:** map of the percentage of converged fits per tile with $\xi > 1.25$. In comparison with the i band map in Fig. 2, it is clear that by applying this cut the overall percentage of successful fits increases dramatically, from $\sim 40\%$ to $> 70\%$ at the borders and up to $\sim 90\%$ in the central areas.

421 $1.25 < \xi \leq 1.5$ and $\xi \geq 1.5$. Values of $\xi < 1$ are unphysical, indicating either noisy photometry, image artefacts or inaccuracies in the PSF model. Each population is represented with a bar coloured by the percentage of converged fits, normalised by the total number of selected objects in each magnitude bin. As expected, we observe lower percentages of converged fits for the objects whose size is comparable to the size of the PSF used by GALFIT to deconvolve their images. Nevertheless, in the range $1 < \xi \leq 1.25$ the completeness is only around 10% lower than at larger sizes. The right panel in Fig. 4 maps the completeness per tile, excluding the galaxy sample whose size is comparable or smaller than the PSF ($\xi < 1.25$). Compared with the i band map in Fig. 2, it shows that by applying the cut in ξ the fitting completeness increases dramatically both at the borders (up to $> 70\%$) and in the central areas (up to $\sim 90\%$), and the discrepancy between these two regions is reduced. In Fig. 3, centre panel, we show the dependence of fitting completeness against PSF FWHM ($i < 21.5$). For the SPREAD_MODEL S/G classifier we see that the completeness at $i < 21.5$ only drops below 80% in the extended tail of the distribution of PSF FWHM (grey histogram).

4.2.3 Image depth

442 There is a clear and expected dependence of the percentage of converged fits on magnitude in both Fig. 2 and Fig. 4. Although stars are less easily excluded at faint magnitudes and the sizes of galaxies are smaller, much of this dependence is likely to be due simply to the difficulty of GALFIT finding a stable minimum in the χ^2 space at low S/N. In the lower panel of Fig. 3 we show how the fitting success rate for $i < 21.5$ galaxies depends on image depth, and hence object S/N. As expected, the completeness falls in shallower regions of the footprint, but the decline is not dramatic for this bright subset and, once again, a high success rate can be maintained by removing only regions corresponding to the tails of the distribution.

4.2.4 Impact of neighbouring sources

454 Finally, we assess the impact of neighbouring sources on the fitting success rate. We reduce the complexity of possible arrangements of neighbours to two metric values: the amount of overlapping area⁵ between a galaxy and its neighbours, and the difference in magnitude between the galaxy and its most overlapping neighbour ($(\text{MAG_AUTO_|C}) - (\text{MAG_AUTO_|MON})$). The dependence of the converged percentage as a function of these two quantities is shown in Fig. 5, in four intervals of S/N for the target object. Each line in the figure is normalized by the population of objects with attempted fits within the same delta-magnitude range. We observe that even at low S/N the fitting success rate is high if all the neighbours present are sufficiently faint. However, in the range $0 < S/N < 25$ the completeness is a steep function of the magnitude difference between target galaxy and its neighbour. At high S/N neither the degree of overlap nor the relative magnitude of a neighbour are important. Note that, our initial selection removes objects that SExtractor determined to have been blended.

4.2.5 Multi-wavelength completeness

473 As shown by the green and red curves in Panel A in Fig. 2, we can recover a relatively high percentage of converged fits for objects brighter than magnitude 21.5 for the g and r filters also. We notice that the g and r bands show a drop in the brightest magnitude range ($\text{GOLD_MAG_AUTO_i} \leq 15.5$). Upon inspection we find that the objects responsible are compact objects with size comparable to the PSF and with a MODEST classification which is close to the threshold of 0.005 in the i -band. In Panels C and D we can see the spatial completeness for the r and g band, respectively. In both cases we reconfirm what we observed for the i band: a poor fitting completeness at the borders of the field, where stellar density is high, as discussed in the previous sub-sections. The g band PSF is typically

⁵ By area, we mean the SExtractor-derived Kron ellipse enlarged by a factor of 1.5

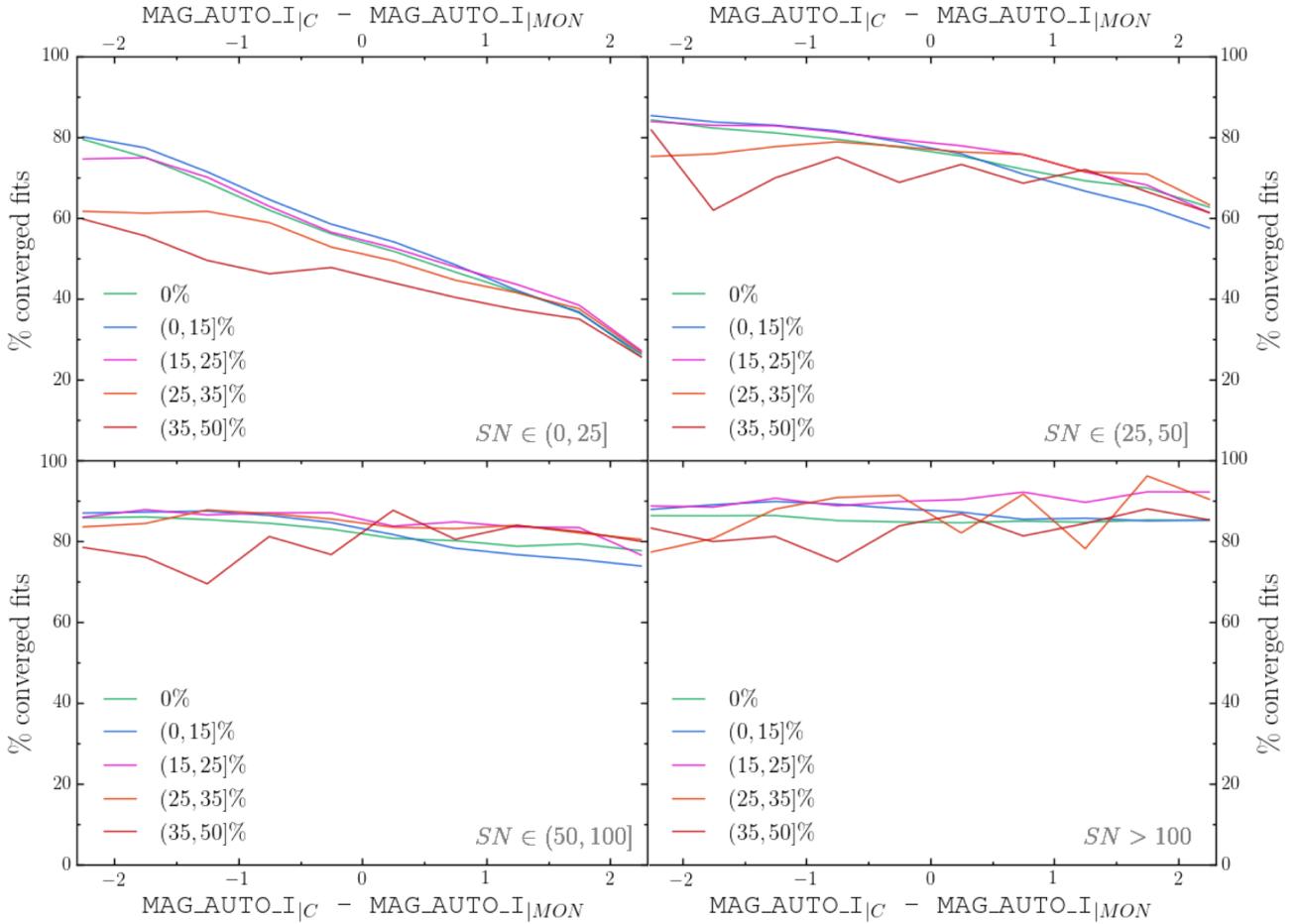


Figure 5. Fitting completeness as a function of the magnitude difference between the target galaxy and its closest neighbour. The relation is shown for different percentages of overlap between the two fitted objects, as reported in the legend. Each line is normalized by the population of objects with attempted fits within the same range in magnitude difference. The analysis is repeated in four signal-to-noise intervals. We observe that the fitting completeness decreases when the closest neighbour is much brighter than the central galaxy, with stronger effects in low signal-to-noise regimes. This effect becomes negligible with increasing signal-to-noise.

485 broader than the r and the i bands, and the images shallower, which
486 are reflected in an overall poorer recovery of converged fits.

487 4.3 Validation

488 We now turn to assessing the accuracy of the parameters recovered
489 from those objects that were successfully fit with GALFIT, begin-
490 ning with simple magnitude and size diagnostics of the population.
491 We then investigate whether there are systematic errors from which
492 GALFIT suffers in recovering the structural parameters of the galax-
493 ies, depending on their magnitude, size, concentration and shape.
494 We investigate this aspect through image simulations (section 2.2)
495 and present the relative calibrations in the next subsection.

496 For this discussion we show the tests performed on the i band,
497 which represents our fiducial filter, starting with a comparison of the
498 total Sérsic magnitude with MAG_AUTO computed by SExtractor.
499 In Fig. 6 we show this comparison for 30,000 randomly-selected
500 objects from the full catalogue. We recover the expected behaviour:
501 objects with Sérsic index ~ 1 have magnitudes consistent with
502 MAG_AUTO, while the Sérsic magnitude is brighter at higher n .
503 MAG_AUTO is known to be biased faint for high-Sérsic n objects,

504 losing as much as 50% of the flux in extreme cases (Graham &
505 Driver 2005).

506 The solid black lines in Fig. 6 delimit the 3σ outliers in magni-
507 tude difference, following an iterative 3-sigma-clipping procedure
508 to find the mean relation and spread (given by the parameters, μ
509 and σ in the figure). The mean relation (red dashed line) is essen-
510 tially flat in magnitude, suggesting that typically the background
511 computed during catalogue extraction and that estimated by GALFIT
512 are consistent. At faint magnitudes, however, there is a population
513 of outliers with magnitude differences that cannot be explained
514 by simple photometric errors, and that also exhibit very high Sérsic
515 indices. We deem these unreliable fits, possibly caused by an
516 unidentified elevated background. Restricting the sample to objects
517 with $S/N > 30$ removes these objects and entirely removes the
518 group with spurious large radii.

519 We then obtain the relation between magnitude and effective
520 radius from the Sérsic profile fits as shown in Fig. 7. Points are
521 colour coded by each object's Sérsic index. Once again, the data
522 match expectations and similar trends reported in the literature,
523 with high Sérsic n objects forming a steep sequence and galaxies
524 with exponential light profiles dominating at fainter magnitudes.

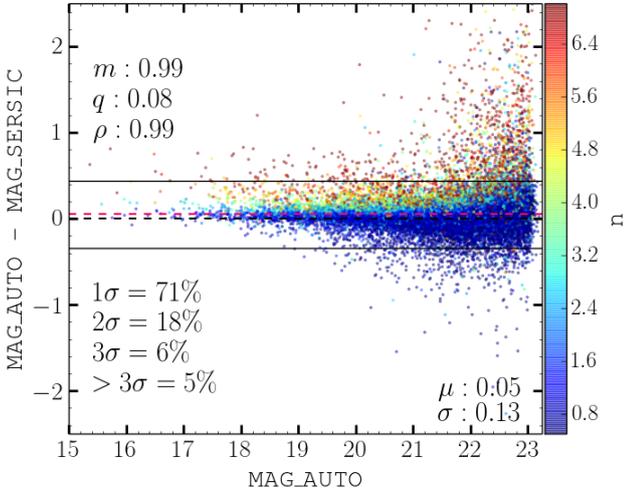


Figure 6. Difference between the input magnitude (MAG_AUTO) from SExtractor and the output magnitude (MAG_SERVIC) recovered through Single Sérsic fits. Results are shown as a function of input magnitude and are colour-coded by Sérsic Index. The two solid black lines delimit the population lying within 3 standard deviations from the mean magnitude difference relation, indicated by the dashed red line. The mean and the spread of the relation, printed in the lower right corner of the Figure, are obtained through a 3σ clipping procedure. The banding in Sérsic index is expected (Graham & Driver 2005) and the vast majority of outliers (which in total number 5% of the sample) are of low S/N objects.

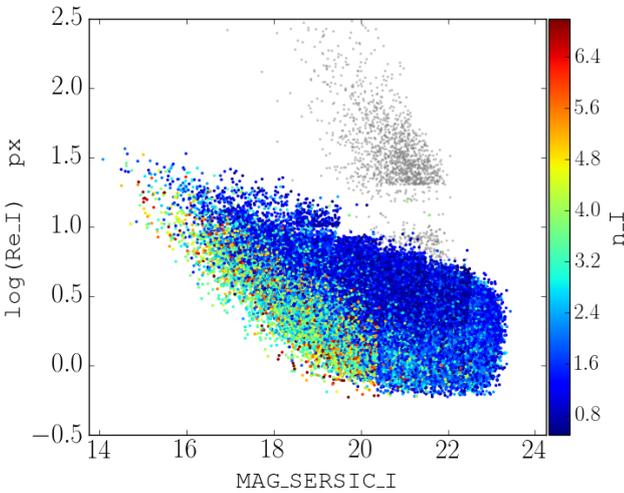


Figure 7. Relation between Sérsic magnitude and effective radius for the i band results. Points are colour-coded by Sérsic Index. outliers are shown in grey.

525 Grey points are sources labelled as outliers during the validation
526 process.

527 4.4 Calibrations

528 In this section we illustrate how we calibrate our measurements. As
529 explained in detail in Section 2.2, we processed and fit the UFIG-
530 BCC simulated data for DES Y1 in the same way we did for our real
531 galaxy sample. Now we can compare the results from the fits with

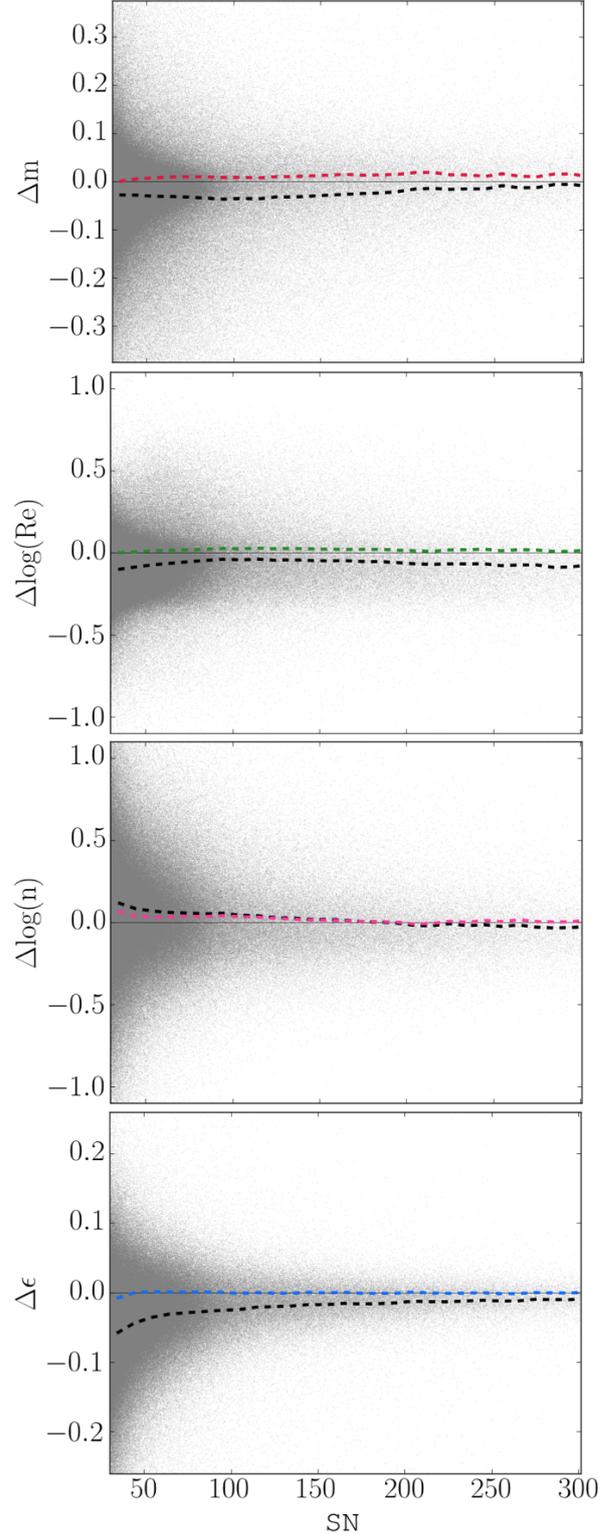


Figure 8. Discrepancies in recovered Sérsic parameters from running GALFIT on the UFIG-BCC image simulations, as a function of signal to noise (S/N). From top to bottom the panels display the results for magnitude, half light radius, Sérsic index and ellipticity. The dashed lines show the discrepancy in bins of S/N, calculated before (black line) and after (coloured line) applying calibration corrections (see section 4.4). The uncertainties depend to first order on the signal to noise, and the scatter is clearly reduced by applying the calibrations. In the calibration map, shown in figure 9, we investigate how the parameters and their uncertainties correlate with each other.

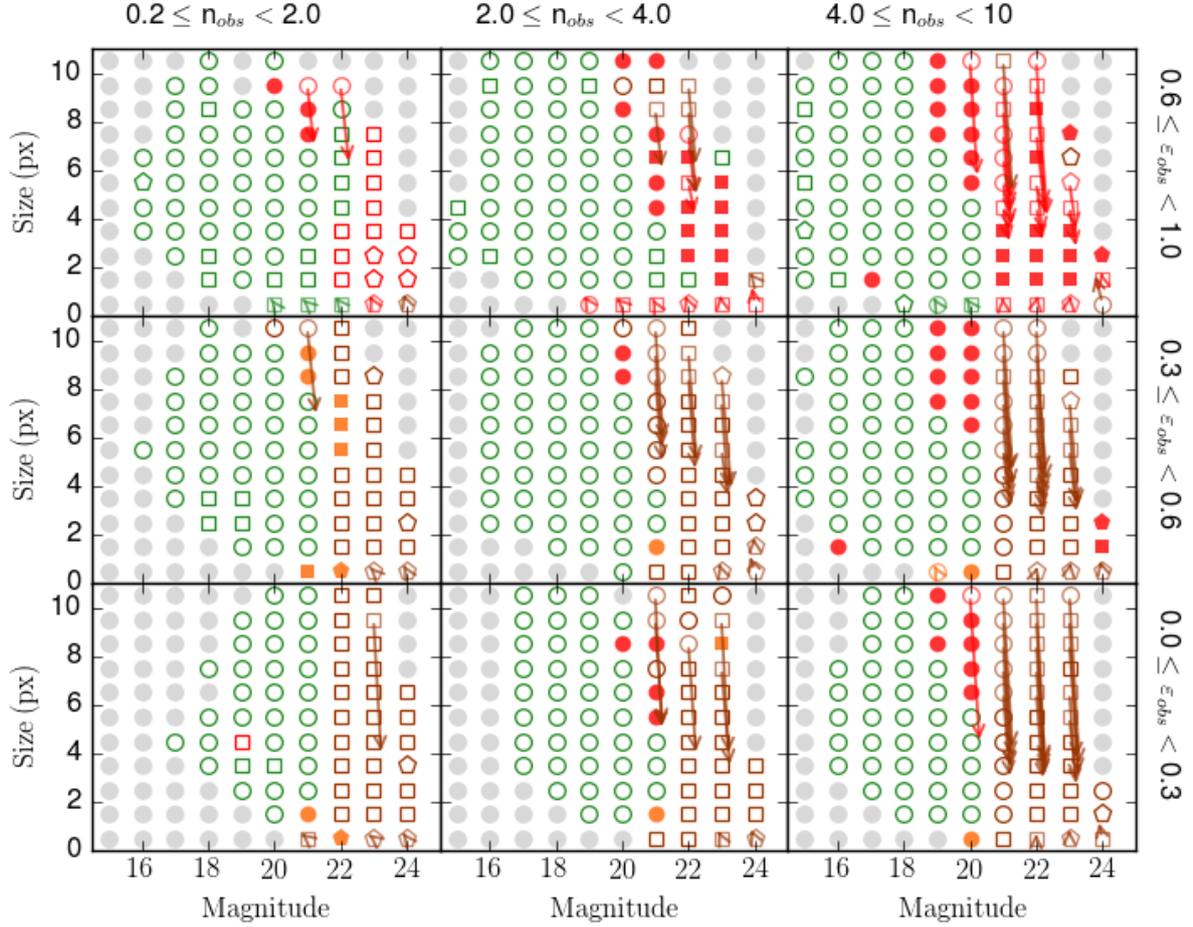


Figure 9. Calibration map for the parametric measurements in the i band, obtained from image simulations as described in Section 2.2. The calibrations are determined in a 4D parameter space, where the correlation of size, magnitude, ellipticity and Sérsic Index between the simulated galaxy and the model is studied. The information is provided using different marker shapes (circles, squares, pentagons, arrows) and colours, as follows. The calibrations are presented in a size-magnitude plane, divided in different cells according to the shown sub-ranges in ellipticity and Sérsic Index. The components of the correction vectors are the magnitude discrepancy η^{mag} and the size discrepancy η^{size} , according to the definitions given in Equations 6 and 7. If these corrections are small ($\eta^{mag} < 0.1 \wedge \eta^{size} < 10\%$) the length of the arrow is set to zero and the cell is identified by a symbol only. Points and arrows are coloured according to the scatter in ellipticity (ϵ) and Sérsic Index (n); a scatter in $\eta^\epsilon > 0.1$ or $\eta^n > 20\%$ is expressed in orange and red, respectively, while the cells presenting a large scatter in both parameters are coloured in brown. The symbol is empty if the GALFIT recovered value is smaller than the model. Different shapes are used referring to the total scatter (w) in the 4D parameter space of the model parameters, defined in Equation 9; the symbol is a pentagon if $w > 1.5$ and a square if $w > 1$, otherwise it is a circle. **These symbols and their meaning are summarised in the legend in Fig. 10**

532 the true morphological parameters used to generate the UFIG-BCC
 533 images. We then calculate the discrepancies between the measured
 534 and true parameters and derive appropriate corrections. We show
 535 the size of these corrections via a set of calibration maps. Symbols
 536 and conventions used in these maps are summarised in the legend
 537 in Fig. 10.

538 4.4.1 Derivation of the corrections

We derive corrections in a 4-dimensional parameter space, including size, magnitude, Sérsic Index and ellipticity. The ensemble of values assumed by each parameter constitutes a vector in the parameter space. We sample each vector with a list of nodes: the magnitude (mag) in the range [14.5,23.5] in steps of 1 magnitude, the size (r) in the interval [0.5,16.5] px in steps of 2 px, the Sérsic Index

(n) in the set [0.2, 2, 4, 10] and the ellipticity (ϵ) in the intervals [0, 0.3, 0.6, 1]. The realization of each combination of these nodes forms an hypervolume which we'll refer to as a *cell*. In each cell falls a certain number of simulated objects with similar structural properties and the corresponding fitting results: so each parameter is represented by a distribution of simulated values and a distribution of measurements. Each distribution in turn has a median value (m^i) and a standard deviation (σ^i), where $i = mag, r, n, \epsilon$, which represent the central value and the dispersion of the population, respectively. To summarise, in each cell the i -th parameter can be expressed as:

$$\hat{i} = \hat{\mu}^i \pm \hat{\sigma}^i \quad (6)$$

for the model and as:

$$i = \mu^i \pm \sigma^i, \quad (7)$$

○	$\eta^\epsilon > 0.1 _{w<1}$	□	$\eta^\epsilon > 0.1 \wedge \eta^n > 20\% _{w>1}$
○	$\eta^n > 20\% _{w<1}$	□	$\eta^{\epsilon,n,mag,size} \sim 0$
○	$\eta^\epsilon > 0.1 \wedge \eta^n > 20\% _{w<1}$	◊	$\eta^\epsilon > 0.1 _{w>1.5}$
○	$\eta^{\epsilon,n,mag,size} \sim 0$	◊	$\eta^n > 20\% _{w>1.5}$
□	$\eta^\epsilon > 0.1 _{w>1}$	◊	$\eta^\epsilon > 0.1 \wedge \eta^n > 20\% _{w>1.5}$
□	$\eta^n > 20\% _{w>1}$	◊	$\eta^{\epsilon,n,mag,size} \sim 0$

Other cases:

symbol+arrow = any previous case $\wedge (\eta^{mag} > 0.1 \wedge \eta^{size} > 10\%)$

filled symbol $\Leftrightarrow \text{GALFIT}_{param} > model_{param}$

Figure 10. Legend for Fig 9 and the calibration maps in Appendix A. It summarises symbols and conventions used in the calibration maps. In the case of the calibration of non-parametric fits, the Sérsic index is replaced with the Concentration parameter.

for the fit, where $i(\hat{i}) = mag, r, n, \epsilon$. For all the objects falling in a given cell we calculate the correction (η^i) in each parameter as the discrepancy between the central values of the distributions:

$$\eta^i = \hat{\mu}^i - \mu^i. \quad (8)$$

We further define a quantity, w , which represents the dispersion of the cell in the 4D parameter space, derived as the quadratic sum of the variances of the model parameters which determine the diagonal of the covariance matrix of the parameter space. It is defined as follows:

$$w = \sqrt{\sum_i \frac{\hat{\sigma}_i^2}{\hat{m}_i^2}}, \quad (9)$$

where $i = mag, r, n, \epsilon$ and $\hat{\sigma}_i^2$ and \hat{m}_i are the variance and median values of the model distributions, respectively. For cells with larger dispersion, we expect the correction vector to be less accurate for a given randomly chosen object.

4.4.2 Calibration maps

In the validation routine we observed that $\sim 99\%$ of converged fits are well recovered in magnitude (η^{mag} of the order of 0.001), and that cutting objects with $S/N < 30$ we remove the clear outliers in size and magnitude. In Figure 8 we show the discrepancies η^i between the intrinsic values and the parametric measurements as a function of signal to noise for magnitude, half-light radius, ellipticity and Sérsic index. The discrepancies relative to size and Sérsic index are shown in logarithmic space to facilitate visualization. In each panel the dashed lines show the discrepancies in bins of signal to noise. We use the uncalibrated sample to calculate the black line, and the same sample after applying the calibrations for the coloured one. It is clear that the uncertainties on the structural parameters increase in low signal to noise regimes, as one might anticipate, and the scatter clearly reduces when applying the corrections. We

observe that GALFIT tends to recover larger sizes and ellipticities, so we pay particular attention to the corrections required for these properties within the multidimensional parameter space.

Figure 9 represents a map of the calibrations that we apply to our measurements, derived from our state-of-the-art image simulations. In using this multidimensional calibration map we are able to account for the correlations between parameters and ensure the corrections are appropriate for a true galaxy sample. The arrows represent the strength of the vector corrections, expressed as the distance between the central values of the size and magnitude distributions of the model sample and the relative measured dataset in each cell. The components of the correction vectors are the magnitude discrepancy η^{mag} on the x axis and the size discrepancy η^{size} on the y axis, according to the definitions given in Equations 6 and 7. If these corrections are small ($\eta^{mag} < 0.1 \wedge \eta^{size} < 10\%$) the length of the arrow is set to zero and only a circle is shown. Apart from the grey circles, which indicate areas with poor statistics, different colours are used to give an indication of the correction applied to ellipticity and Sérsic Index. If $\eta^\epsilon > 0.1$ or $\eta^n > 20\%$, the symbol is coloured in orange and red, respectively. If the correction is large in both cases, then it is coloured in brown. The symbol is empty if the GALFIT recovered value is smaller than the model. The symbols are shaped according to the total scatter (w) in the 4D parameter space of the model parameters, defined in Equation 9; we use a pentagon if $w > 1.5$ and a square if $w > 1$, otherwise the symbol is a circle. Figure 9 reports the vector corrections for the i band; corrections for the g and r filters are shown in the Appendix A.

We observe that the strength of the corrections and their positions are compatible with the findings we discussed previously in the validation section. In that section we noted that in any range of shape and Sérsic index the uncalibrated measurements of the sub-populations of galaxies at the faintest magnitude range present overestimated half light radii and Sérsic Indices. In the calibration map they are assigned with larger vector corrections in size, which calibrate the measurements towards smaller values. If the correction

in size is small, then we observe that a calibration in Sérsic Index is applied, where the recovered value was larger than the model parameter. The same observations are valid also for the other two filters (shown in Appendix A). The fact that the measurements and their associated corrections are similar across photometric bands indicates that our final set of calibrated results are robust to the survey characteristics, such as overall PSF size and noise level, that vary between bands. Furthermore, the vast majority of cells across all three calibration maps show little corrections, suggesting that our converged fits are in general reliable and represent the light profiles well.

5 NON PARAMETRIC FITS

5.1 ZEST+ Setup

ZEST+ is a C++ software application which uses a non-parametric approach to quantify galaxy structure and perform morphological classification. It is based on the ZEST algorithm by Scarlata et al. 2007a,b, which saw a first application in Cameron et al. 2010. Compared with its predecessor, ZEST+ has increased execution speed. The software architecture consists of two main modules: *Preprocessing* and *Characterization*. The former performs image cleaning, main object centring and segmentation, the latter calculates structure and substructure morphological coefficients.

5.1.1 Preprocessing

In this module the algorithm uses the stamps and the input catalogue provided by the SAND routine. The input catalogue includes the coordinates and the geometrical parameters of the target galaxy and its neighbours in order to remove nearby objects, subtract the background, determine the centre of the galaxy and measure its Petrosian radius.

The Petrosian radius is defined as the location where the ratio of flux intensity at that radius, $I(R)$, to the mean intensity within the radius, $\langle I(< R) \rangle$, reaches some value, denoted by $\eta(R)$ (Petrosian 1976):

$$\eta = \frac{I(R)}{\langle I(R) \rangle}. \quad (10)$$

For this work the Petrosian radius corresponds to the location where $\eta(R) = 0.2$. The Petrosian ellipse associated with the object contains the pixels which are used in the *Characterization* module to calculate the morphological coefficients of the central galaxy.

5.1.2 Characterization

The measurements provided by ZEST+ are galaxy concentration (C), asymmetry (A), clumpiness or smoothness (S) and Gini (G) and M_{20} coefficients. This set of parameters, which we refer to as to the *CASGM system*, quantifies the galaxy light distribution and is widely used in studies which correlate the galaxy structure to other parameters, such as colour and peculiar features indicating mergers or galaxy interactions (see for example Conselice et al. 2000, Conselice 2003, Lotz et al. 2004 and Zamojski et al. 2007); other similar quantities have been recently introduced by Freeman et al. (2013).

The *concentration* of light, first introduced in Bershadly et al. 2000

and Conselice 2003, expresses how much light is in the centre of a galaxy as opposed to its outer parts; it is defined as

$$C = 5 \log \left(\frac{r_{80}}{r_{20}} \right), \quad (11)$$

where r_{80} and r_{20} are the elliptical radii enclosing, respectively, the 20% and 80% of the flux contained within the Petrosian ellipse of the object. ZEST+ outputs three different values of concentration, C , C_{ext} and C_{circ} . The first parameter is calculated using the total flux measured within the Petrosian ellipse, the second using the flux given as input by the user within the same ellipse and the third one using the Petrosian flux within a circular aperture. For this work we refer to C as the concentration.

The *asymmetry* is an indicator of what fraction of the light in a galaxy is in non-symmetric components. Introduced in Schade et al. 1995 first, and then in Abraham et al. 1996 and Conselice 1997 independently, asymmetry is determined by rotating individual galaxy images by 180° about their centres and self-subtracting these from the original galaxy images. This procedure is applied after the *Preprocessing* module, where the background is $\kappa\sigma$ -clipped and subtracted. The value of pixel (i, j) in the subtracted image is calculated as:

$$\Delta I(i, j) = I(i, j) - I_{180}(i, j) = I(i, j) - I(2i_c - i, 2j_c - j), \quad (12)$$

where I_{180} is the rotated image and (i_c, j_c) are coordinates of the centre of the galaxy.

To take into account the asymmetry of the background, ZEST+ follows the same method, as in Zamojski et al. 2007, working with smoothed images of the galaxies and their rotated version. The asymmetry of the original image is defined as

$$A_0 = \frac{1}{2} \frac{\sum_{i,j} |I(i, j) - I_{180}(i, j)|}{\sum_{i,j} |I(i, j)|}, \quad (13)$$

where $I(i, j)$ and $I_{180}(i, j)$ express the intensity of the flux at the pixel (i, j) in the original and rotated image, respectively. Similarly we define the asymmetry of the smoothed image:

$$A_{0,S} = \frac{1}{2} \frac{\sum_{i,j} |I^S(i, j) - I_{180}^S(i, j)|}{\sum_{i,j} |I^S(i, j)|}. \quad (14)$$

Assuming that the intrinsic asymmetry of the light does not change in the smoothed version, we consider that the difference between the two values of asymmetry is due to the background. Smoothing reduces the standard deviation of the background by a factor $\sqrt{5}$ with respect to its un-smoothed version. The combination of A_0 and $A_{0,S}$ then gives the final asymmetry value:

$$A = A_0 - \frac{A_0 - A_{0,S}}{1 - 1/\sqrt{5}}, \quad (15)$$

where the subtracted term corresponds to the background correction factor.

The *clumpiness* or *smoothness* parameter, introduced in Conselice 2003, describes the fraction of light which is contained in clumpy distributions. Clumpy galaxies show a large amount of light at high spatial frequencies, and smooth systems at low frequencies. This parameter is therefore useful to catch patches in the galaxy light which reveal star-forming regions and other fine structure. ZEST+ calculates the clumpiness by subtracting a smoothed image, $I_S(i, j)$, from the original, $I(i, j)$, and then quantifying the residual image, $I_\Delta(i, j)$. The smoothed image is obtained by convolving the original image with a Gaussian filter of FWHM equal to 0.25

times the Petrosian radius calculated during the *Preprocessing* module. In $I_{\Delta}(i, j)$ the clumpy regions are quantified from the pixels with intensity higher than $k = 2.5$ times the background standard deviation in the residual image σ_{Δ} . These pixels are then used to calculate the clumpiness of the galaxy:

$$S = \frac{\sum_{i,j} I_{\Delta}(i, j)}{\sum_{i,j} |I(i, j)|_{I_{\Delta}(i, j) > k\sigma_{\Delta}}} \quad (16)$$

Similarly, the *Gini* coefficient quantifies how uniformly the flux of an object is distributed among its pixels. A Gini coefficient $G = 1$ indicates that all the light is in one pixel, while $G = 0$ means that every pixel has an equal share. To calculate *Gini* ZEST+ uses the definition by Lotz et al. (2004, 2008a,b):

$$G = \frac{1}{\hat{I}n(n-1)} \sum_n^n (2i - n - 1)\hat{I}_i, \quad (17)$$

where \hat{I} is the mean flux of the galaxy pixels.

The M_{20} coefficient is similar to the concentration C in that its value indicates the degree to which light is concentrated in an image; however a high light concentration (denoted by a very negative value of M_{20}) doesn't imply a central light concentration. For this reason it is useful in describing the spatial distribution of bright substructures within the galaxy, such as spiral arms, bars or bright nuclei. The computation of this parameter requires first that the pixels within the Petrosian ellipse of the galaxy are ordered by flux; then the 20% brightest pixels are selected and for each pixel i the second-order moments are calculated:

$$E_i = I_i[(x_i - x_c)^2 + (y_i - y_c)^2], \quad (18)$$

where I_i is the flux in the i -th pixel, (x_i, y_i) the coordinates of the pixel and (x_c, y_c) the coordinates of the centre of the Petrosian ellipse. The sum of these moments is $E = \sum_i^{N_{20}} E_i$, where N_{20} is the multiplicity of the 20% brightest selected pixels. Given E_{tot} as the sum of the second order moments of all the pixels in the ellipse, we finally calculate M_{20} as:

$$M_{20} = \text{Log} \frac{E}{E_{tot}}. \quad (19)$$

5.2 Completeness

The measurements of Gini, M_{20} , Concentration, Asymmetry and Clumpiness are matched with diagnostic flags which inform the user whether errors occurred during the cleaning step of the process or in the calculation of the coefficients. To be more precise, the flag **Error** (we label it in our catalogue as **ERRORFLAG**) indicates whether a problem occurred while processing an object: if it is non-zero, it traces an error encountered during the calculation of the structural parameters, and flags the measurements as not reliable. The *contamination flag* informs the user whether the cleaning process was unsuccessful due to the presence of a neighbour covering the centre of the galaxy; in this case the program outputs *contamination flag* = -2. Therefore in this test we considered as converged fits the measurements with **ERRORFLAG** = 0 \wedge *contamination flag* \neq -2. Then we define the fitting completeness as we did for the parametric fits, following Equation 3.

The results for the g , r and i bands are shown in Figure 11. With the cut in **ERRORFLAG** and *contamination flag* we discard a

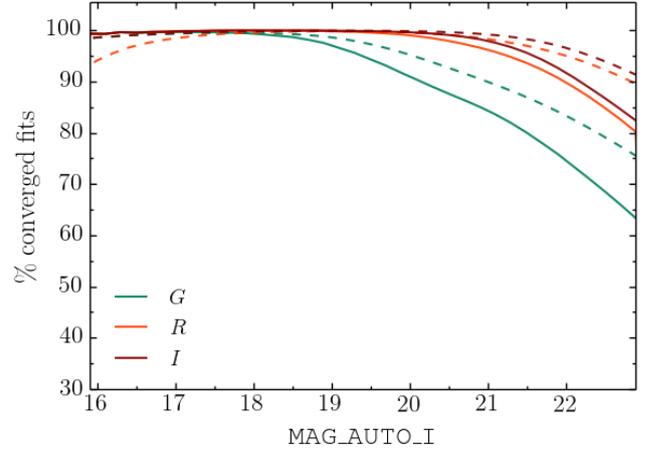


Figure 11. Fitting completeness of non-parametric converged fits in the g , r and i bands, expressed in terms of the percentage of converged fits in bins of 0.2 magnitude, normalised on the total number of selected objects in that magnitude bin. By *converged fits* we refer in this case to the objects flagged by ZEST+ as fits without errors, either during the cleaning process or the characterization routine, as described in more detail in Section 5.2. Magnitude-limited completeness is represented by the dashed lines. We obtain almost full recovery in the i and r filters up to $i \sim 21$, losing only a few saturated objects.

total of $\sim 10\%$ of objects. We observe some fluctuations at the brightest end, where we find cases of large bright galaxies whose Petrosian ellipses were underestimated or cases with saturated objects, and at the faintest end, where it is more common to have higher noise contamination within the Petrosian ellipse. The overall number of successful fits is more than $\sim 90\%$ in the i and r filters and $\sim 80\%$ in the g band. The dashed lines show magnitude-limited, rather than differential, completeness.

5.3 Validation

By way of a simple internal validation, we show in Figure 12 the uncalibrated measurements from ZEST+ and the relationships between them. In particular we focus on the Gini- M_{20} relation, studied as a function of other morphological parameters: Concentration (C), Clumpiness (S), Asymmetry (A) and Ellipticity (ϵ). We observe that even though the measurements are still un-calibrated, we can easily recover the expected trends with very few outliers. As an example consider the first panel, where the Gini- M_{20} relation is colour-coded by the Concentration. The objects with low M_{20} values present high concentration of light; from the figure we observe that in the Gini- M_{20} plane these objects tend to have larger values of Gini, which means that the light is not uniformly distributed; if we now add the third parameter, we notice that the Concentration of these objects lies in its highest range: this explains that the light of these galaxies is very concentrated, and locates it at the centre of the galaxy. From panels c, b and d we also add the expected information that these objects are symmetric, lack clumpy regions and are mostly rounded. These observations are also confirmed by visual inspection of image stamps.

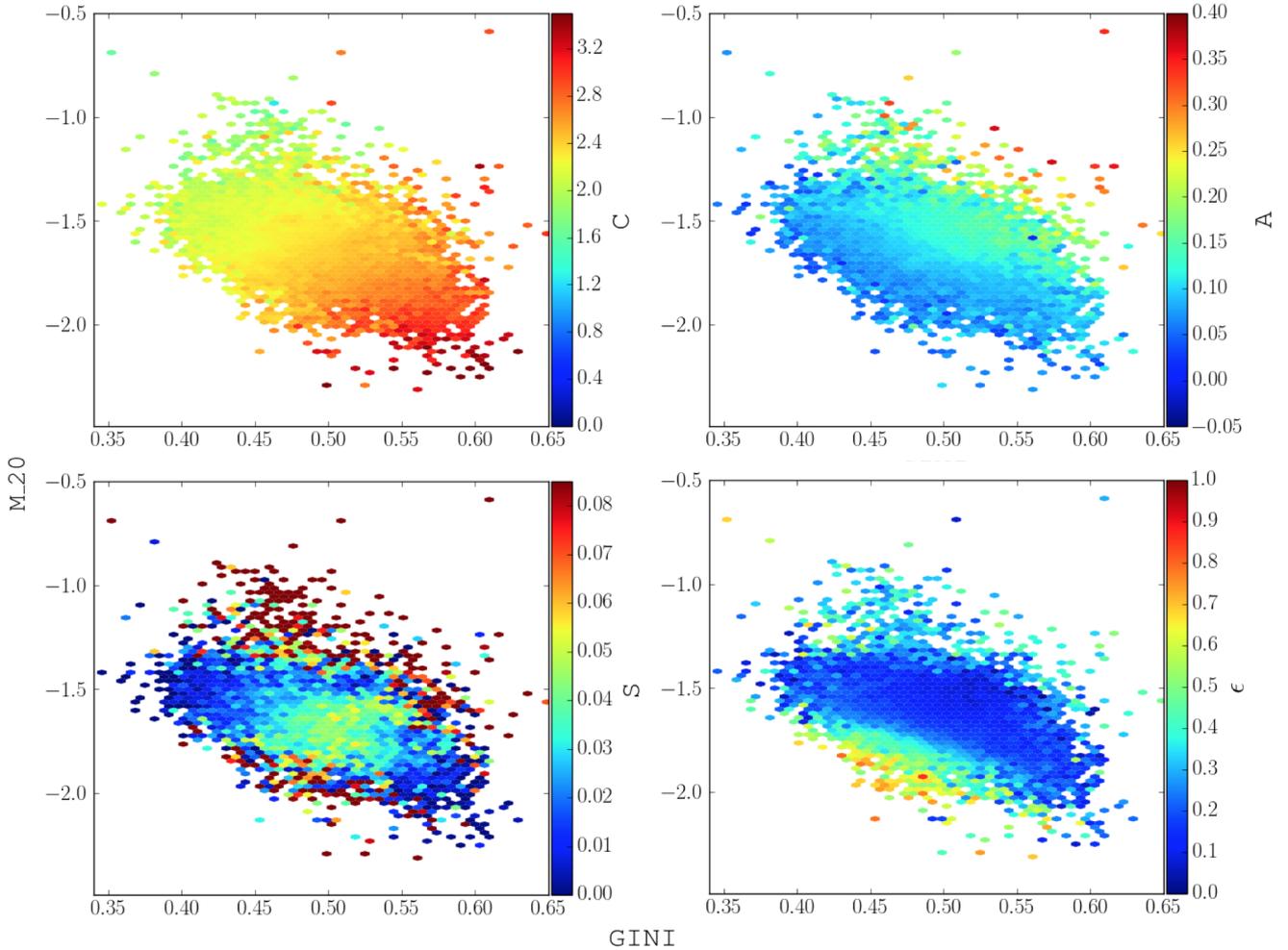


Figure 12. Gini-M20 relation shown as a function of Concentration C (panel A), Clumpiness S (panel B), Asymmetry A (panel C) and Ellipticity ϵ (panel D). The expected trends for the relations and their gradients are recovered, as discussed in more detail in Section 5.3.

5.4 Calibrations and diagnostics of the corrected results

In order to apply corrections to the non-parametric measurements, which are crucial in accounting for the impact of the PSF, we adopt the same approach used for the parametric fits: we consider the images from the *UFig-BCC* release for DES Y1 and treat them as if they were real data, as explained in detail in Section 2.2. We then derive calibration maps exactly as described in Section 4.4.1, determining the correction for each parameter of interest as the discrepancy between the central values of the model and the fitting results distributions in each cell. The equations 6, 7, 8 and 9 are valid also in this context, with the exception that the Sérsic Index, n , is now substituted by the Concentration of light, C .

In order to derive correction vectors, we first compute ZEST+ output parameters for the simulated galaxies *before* noise and PSF convolution are applied. We use GALFIT to produce noise and PSF-free image stamps based on the UFig model parameters and run ZEST+ on them. In this way we construct the truth table of values with which to derive calibration vectors. Figure 13 shows the correction map for the i band; the other two filters, g and r , are presented in Appendix A. Also for non-parametric fits we adopt the same convention of colours and shapes as in Figure 9. The length

of the arrows is a visual representation of the strength of the vector correction: their x and y components are the discrepancies between the central values of the model distribution and the fitted dataset in each 4-dimensional cell, projected on the size-magnitude plane. When the correction is small ($\eta^{mag} < 0.1 \wedge \eta^{size} < 10\%$) a symbol in place of the arrow is shown. Apart from the grey circles, which indicate areas with poor statistics, the colour legend reflects the size of the calibration applied to ellipticity and Concentration. If the scatter in ellipticity or Concentration is large ($\eta^\epsilon > 0.1$ or $\eta^C > 20\%$), then the symbol is coloured in orange or red, respectively. If this condition applies to both parameters simultaneously, it is coloured in brown. If the recovered value underestimates the model input, the symbol is empty. Different shapes are used according to the dispersion w of the 4-dimensional parameter space, calculated considering its covariance matrix, as expressed in Equation 9. Symbols are pentagons when $w > 1.5$, squares if $w > 1$ and circles otherwise. We observe that the majority of red cells, where a larger correction in Concentration is required, have an empty symbol: this tells us that ZEST+ tends to recover underestimated values of concentration. This behaviour is entirely expected, due to the fact that ZEST+ cannot account the PSF in computing results. We demonstrate this aspect more explicitly in Figure 15, which shows

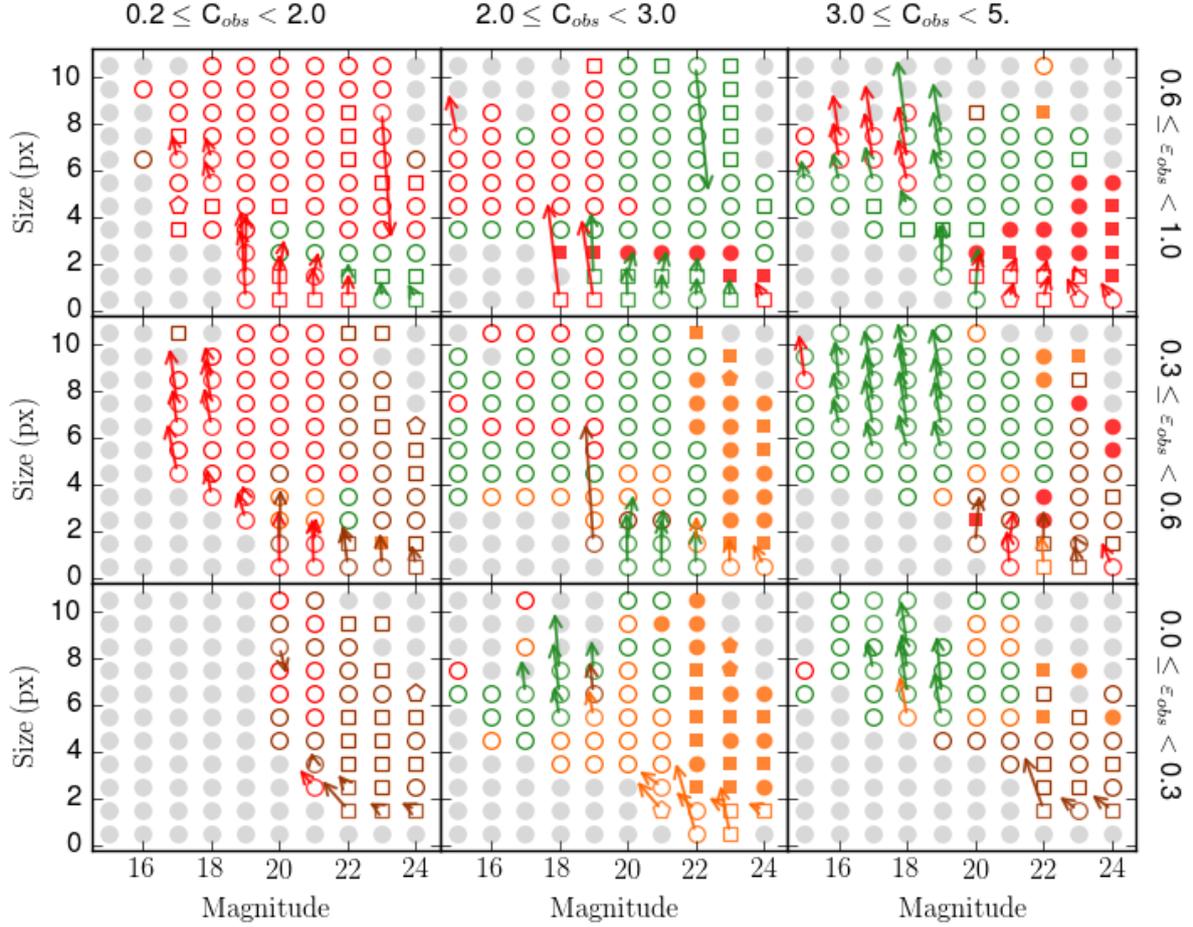


Figure 13. Calibration map for the non-parametric measurements in the i band, obtained through the simulation routine described in Section 5.4. The calibrations are determined in a 4D parameter space, where the correlation of size, magnitude, ellipticity and Concentration between the measured values and the model parameters is studied. The information in the map is displayed using different symbols and colours with the same GALFIT adopted for the parametric fits. They calibrations are presented in a size-magnitude plane, divided in different cells according to the shown sub-ranges in ellipticity and Concentration. The components of the correction vectors are the magnitude discrepancy η^{mag} on the x axis and the size discrepancy η^{size} on the y axis, according to the definitions given in Equations 6 and 7. If these corrections are small ($\eta^{mag} < 0.1 \wedge \eta^{size} < 10\%$) the length of the arrow is set to zero and only a symbol identifies them. If the scatter in ellipticity (ϵ) or Concentration (C) is large ($\eta^\epsilon > 0.1$ and $\eta^C > 20\%$, respectively), then the symbol is coloured in orange or red, respectively. If the calibration is large in both parameters, it is coloured in brown. The symbol is empty if the ZEST+ recovered value is smaller than the model. Different shapes are used referring to the total scatter (w) in the 4D parameter space of the model parameters, defined in Equation 9; the symbol is a pentagon if $w > 1.5$ and a square if $w > 1$, otherwise it is a circle.

713 the relation between the Sérsic Index and the Concentration before
714 (black contours) and after (magenta) applying the corrections: the
715 *flattening* effect we observe in the uncalibrated population of Con-
716 centration values reflects exactly what we observe in the calibration
717 map. It is then shifted towards more realistic values through the
718 corrections. This test shows that using calibrated values from both
719 parametric and non-parametric approaches to quantifying galaxy
720 structure allows us to use the advantages of both methods and provide
721 a firmer grip on the characteristics of the galaxy population.
722 We will exploit the strength of our dual-method, multi-band morpho-
723 logic catalogue in a series of future papers.

724 6 SCIENCE-READY CUTS

725 We finish by summarising the overall selection function of the
726 galaxy sample and detail a set of simple cuts that could form the ba-
727 sis of a sample for scientific analysis. We exclude from consideration
728 objects that meet any one of the following criteria:

- 729 • SExtractor FLAGS > 0
- 730 • CLASS_STAR > 0.9
- 731 • MAG_AUTO_I > 23
- 732 • FLUX_RADIUS ≤ 0
- 733 • KRON_RADIUS ≤ 0
- 734 • FLAGS_BADREGION > 0
- 735 • Objects with a neighbour that overlaps 50% or more of its
736 expanded Kron ellipse. The relevant column in the catalogue for
737 this criterion is MAX_OVERLAP_PERC.

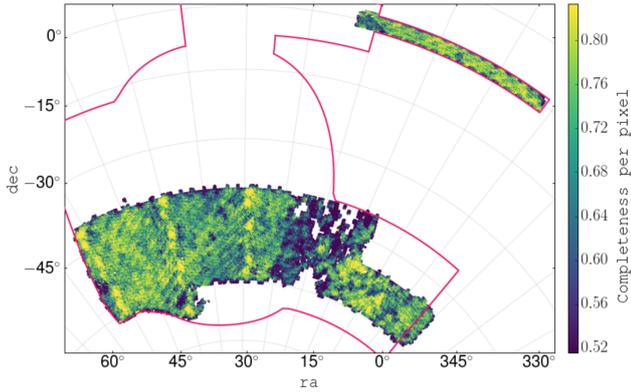


Figure 14. Healpix map of the ratio between two galaxy samples. We apply to the Y1A1 data the sample-selection cuts to obtain the first sample, and then apply the science-ready cuts to it in order to get the second one. The ratio gives the completeness per pixel of the science-ready sample.

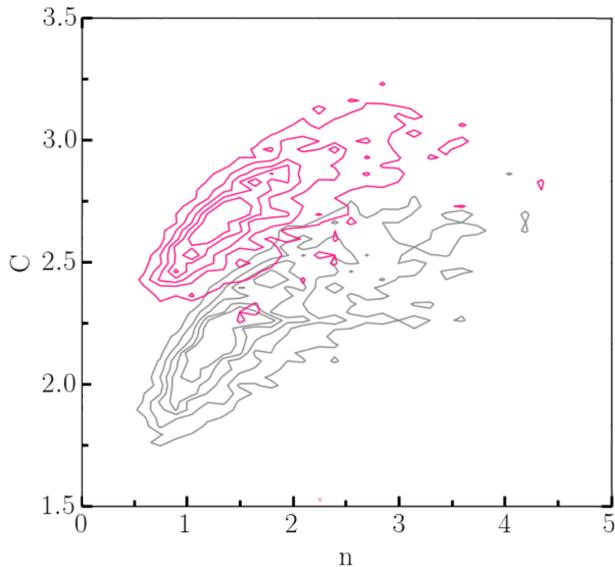


Figure 15. Sérsic Index-Concentration relation before (grey) and after (magenta) applying the calibrations. The *flattening effect* present in the uncalibrated measurements is due to PSF effects which is solved via the calibrations.

- Objects that have unrecoverable errors in the SExtractor output of their neighbouring objects (if any).

This initial sample comprises 45 million objects over 1800 square degrees that is 80% complete in Sérsic measurements up to magnitude 21.5.

To prepare a high completeness science-ready galaxy sample, we suggest the following initial cuts. Science problems requiring higher completeness and/or greater uniformity across the footprint will require additional cuts, dependent on the goals. In some circumstances fainter galaxies could also be included in the sample.

- $\text{MAG_AUTO_I} \leq 21.5$
- $S/N > 30$
- $\text{SPREAD_MODEL} + 1.67 \times \text{SPREADERR_MODEL} > 0.005$

For the i-band catalogue, these cuts produce a sample of 12 million galaxies that is 90% complete in Sérsic measurements and 99% complete in non-parametric measurements.

In Fig. 14 we show a ratio of two healpix maps realised with two samples. We first applied the cuts used for the sample selection, with an additional cut in $\text{MAG_AUTO} < 21.5$. We chose this threshold according to the analysis of the completeness discussed in Section 4.2. Then we select from this sample all the objects with pass the set of science-ready cuts we proposed above. The map shows the completeness per pixel, which is overall uniform. It also guides the catalogue users to possibly select specific areas for future analyses.

7 CONCLUSIONS

We have presented the process of preparing, producing and assembling the largest structural and morphological galaxy catalogue to date, comprising 45 million objects over 1800 square degrees, which are taken from the first year of the Dark Energy Survey observations (DES Y1). We adopted both parametric and non-parametric approaches, using GALFIT and ZEST+. In order to optimize their performance according to the characteristics of our sample, in particular in those cases where the galaxy we want to fit has one or more close neighbours, we developed a neighbour-classifier algorithm as part of a pre-fitting pipeline (Section 3.2) which automatically prepares the postage stamps and all the settings required to simultaneously fit the objects in the presence of overlapping isophotes. We stress the importance of this step because a precise treatment of the size of the stamps and the neighbouring objects allows the recovery of more accurate measurements.

In Section 4.2 we presented the fitting completeness of the parametric fits in the g , r and i filters as a function of object magnitude. Using a tile-by-tile analysis, we show that the highest percentages of non-converged fits are localised at the West and East borders of the footprint, where there is a high stellar density due to the vicinity of the Large Magellanic Cloud. After applying star-galaxy separation based on a linear combination of the parameter SPREAD_MODEL and its uncertainty, we find that the fitting efficiency remains high ($> 80\%$) up to magnitude < 22 for the i and r band, and magnitude < 21 for the g band. We also studied the subsequent fitting completeness in relation to survey data characteristics that are expected to impact the performance of GALFIT: stellar density, PSF FWHM and image depth. We conclude that at relatively bright magnitudes ($i < 21.5$) the completeness has a relatively weak dependence on these quantities, and high completeness can be maintained without much loss of survey area.

In Section 4.3 we analysed the properties of the converged fits, isolating a small fraction ($< 5\%$) of outliers in magnitude recovery, and a branch of objects with high Sérsic indices and large radii that we believe to be spurious. Removing low S/N galaxies efficiently cleans the sample of these populations. Following this basic validation, we calibrate the Sérsic measurements using state-of-the-art UFIG image simulations, deriving correction vectors via the comparison of input model parameters and the resulting fits by GALFIT. In Section 5 we repeated the above mentioned diagnostics for the non-parametric fits, benefiting from the internal diagnostic flags provided by ZEST+ itself in order to quantify the quality of the image and so the reliability of the measurements. For the non-parametric dataset we adopted the same method to derive the calibrations described in Section 2.2, finding that corrections are stronger for low signal to noise galaxies, similar to the parametric case. In particular, we highlight the calibration of galaxy concentration, which is adversely affected due to

810 fact that ZEST+ cannot account for the PSF.
 811 Finally, we summarised the selection function and a recommended
 812 set of cuts to form a basic science sample. Our catalogue represents
 813 a valuable instrument to explore the properties and the evolutionary
 814 paths of galaxies in the DES Y1 survey volume, which will be used
 815 in a series of forthcoming publications.

816 **ACKNOWLEDGEMENTS**

817 FT would like to thank Sandro Tacchella for his useful suggestions
 818 about fitting calibration and fruitful discussions.

820 Funding for the DES Projects has been provided by the U.S. De-
 821 partment of Energy, the U.S. National Science Foundation, the Min-
 822 istry of Science and Education of Spain, the Science and Technol-
 823 ogy Facilities Council of the United Kingdom, the Higher Education
 824 Funding Council for England, the National Center for Supercomput-
 825 ing Applications at the University of Illinois at Urbana-Champaign,
 826 the Kavli Institute of Cosmological Physics at the University of
 827 Chicago, the Center for Cosmology and Astro-Particle Physics at
 828 the Ohio State University, the Mitchell Institute for Fundamental
 829 Physics and Astronomy at Texas A&M University, Financiadora
 830 de Estudos e Projetos, Fundação Carlos Chagas Filho de Amparo à
 831 Pesquisa do Estado do Rio de Janeiro, Conselho Nacional de Desen-
 832 volvimento Científico e Tecnológico and the Ministério da Ciência,
 833 Tecnologia e Inovação, the Deutsche Forschungsgemeinschaft and
 834 the Collaborating Institutions in the Dark Energy Survey.

835 The Collaborating Institutions are Argonne National Labora-
 836 tory, the University of California at Santa Cruz, the University of
 837 Cambridge, Centro de Investigaciones Energéticas, Medioambien-
 838 tales y Tecnológicas-Madrid, the University of Chicago, Univer-
 839 sity College London, the DES-Brazil Consortium, the University
 840 of Edinburgh, the Eidgenössische Technische Hochschule (ETH)
 841 Zürich, Fermi National Accelerator Laboratory, the University of
 842 Illinois at Urbana-Champaign, the Institut de Ciències de l’Espai
 843 (IEEC/CSIC), the Institut de Física d’Altes Energies, Lawrence
 844 Berkeley National Laboratory, the Ludwig-Maximilians Universität
 845 München and the associated Excellence Cluster Universe, the Uni-
 846 versity of Michigan, the National Optical Astronomy Observatory,
 847 the University of Nottingham, The Ohio State University, the Uni-
 848 versity of Pennsylvania, the University of Portsmouth, SLAC Na-
 849 tional Accelerator Laboratory, Stanford University, the University
 850 of Sussex, Texas A&M University, and the OzDES Membership
 851 Consortium.

852 Based in part on observations at Cerro Tololo Inter-American
 853 Observatory, National Optical Astronomy Observatory, which is op-
 854 erated by the Association of Universities for Research in Astronomy
 855 (AURA) under a cooperative agreement with the National Science
 856 Foundation.

857 The DES data management system is supported by the Na-
 858 tional Science Foundation under Grant Numbers AST-1138766
 859 and AST-1536171. The DES participants from Spanish institu-
 860 tions are partially supported by MINECO under grants AYA2015-
 861 71825, ESP2015-66861, FPA2015-68048, SEV-2016-0588, SEV-
 862 2016-0597, and MDM-2015-0509, some of which include ERDF
 863 funds from the European Union. IFAE is partially funded by the
 864 CERCA program of the Generalitat de Catalunya. Research leading
 865 to these results has received funding from the European Research
 866 Council under the European Union’s Seventh Framework Program
 867 (FP7/2007-2013) including ERC grant agreements 240672, 291329,
 868 and 306478. We acknowledge support from the Australian Research

869 Council Centre of Excellence for All-sky Astrophysics (CAAS-
 870 TRO), through project number CE110001020, and the Brazilian In-
 871 stituto Nacional de Ciência e Tecnologia (INCT) e-Universe (CNPq
 872 grant 465376/2014-2).

873 This manuscript has been authored by Fermi Research Al-
 874 liance, LLC under Contract No. DE-AC02-07CH11359 with the
 875 U.S. Department of Energy, Office of Science, Office of High En-
 876 ergy Physics. The United States Government retains and the pub-
 877 lisher, by accepting the article for publication, acknowledges that
 878 the United States Government retains a non-exclusive, paid-up, ir-
 879 revocable, world-wide license to publish or reproduce the published
 880 form of this manuscript, or allow others to do so, for United States
 881 Government purposes.

882 **REFERENCES**

883 Abbott T., et al., 2016, *MNRAS*, **460**, 1270
 884 Abraham R. G., van den Bergh S., Glazebrook K., Ellis R. S., Santiago
 885 B. X., Surma P., Griffiths R. E., 1996, *ApJS*, **107**, 1
 886 Abraham R. G., van den Bergh S., Nair P., 2003, *ApJ*, **588**, 218
 887 Bamford S. P., et al., 2009, *MNRAS*, **393**, 1324
 888 Banerji M., et al., 2010, *MNRAS*, **406**, 342
 889 Barden M., Häubler B., Peng C. Y., McIntosh D. H., Guo Y., 2012, *MNRAS*,
 890 **422**, 449
 891 Barnes J. E., 1988, *ApJ*, **331**, 699
 892 Bell E. F., et al., 2012, *ApJ*, **753**, 167
 893 Bergé J., Gamper L., Réfrégier A., Amara A., 2013, *Astronomy and Com-*
 894 *puting*, **1**, 23
 895 Bernstein R. A., Freedman W. L., Madore B. F., 2002, *ApJ*, **571**, 56
 896 Bershady M. A., Jangren A., Conselice C. J., 2000, *AJ*, **119**, 2645
 897 Bertin E., 2011, in Evans I. N., Accomazzi A., Mink D. J., Rots A. H.,
 898 eds, *Astronomical Society of the Pacific Conference Series Vol. 442*,
 899 *Astronomical Data Analysis Software and Systems XX*. p. 435
 900 Bertin E., 2013, PSFEX: Point Spread Function Extractor, *Astrophysics*
 901 *Source Code Library* (ascl:1301.001)
 902 Bertin E., Arnouts S., 1996, *A&AS*, **117**, 393
 903 Blanton M. R., et al., 2003, *ApJ*, **594**, 186
 904 Bournaud F., Elmegreen B. G., Elmegreen D. M., 2007, *ApJ*, **670**, 237
 905 Busha M. T., Wechsler R. H., Becker M. R., Erickson B., Evrard A. E., 2013,
 906 in *American Astronomical Society Meeting Abstracts #221*. p. 341.07
 907 Cameron E., et al., 2010, *MNRAS*, **409**, 346
 908 Combes F., Sanders R. H., 1981, *A&A*, **96**, 164
 909 Conselice C. J., 1997, *PASP*, **109**, 1251
 910 Conselice C. J., 2003, *ApJS*, **147**, 1
 911 Conselice C. J., 2008, in Knäpen J. H., Mahoney T. J., Vazdekis A., eds, *As-*
 912 *tronomical Society of the Pacific Conference Series Vol. 390*, *Pathways*
 913 *Through an Eclectic Universe*. p. 403 ([arXiv:0706.3482](https://arxiv.org/abs/0706.3482))
 914 Conselice C. J., Bershady M. A., Jangren A., 2000, *ApJ*, **529**, 886
 915 Conselice C. J., Rajgor S., Myers R., 2008, *MNRAS*, **386**, 909
 916 DES Collaboration 2005, *ArXiv Astrophysics e-prints*,
 917 Dekel A., et al., 2009a, *Nature*, **457**, 451
 918 Dekel A., Sari R., Ceverino D., 2009b, *ApJ*, **703**, 785
 919 Desai S., et al., 2012, *ApJ*, **757**, 83
 920 Diehl H. T., et al., 2014, in *Observatory Operations: Strategies, Processes,*
 921 *and Systems V*. p. 91490V, [doi:10.1117/12.2056982](https://doi.org/10.1117/12.2056982)
 922 Dieleman S., Willett K. W., Dambre J., 2015, *MNRAS*, **450**, 1441
 923 Dressler A., 1980, *ApJ*, **236**, 351
 924 Drlica-Wagner A., et al., 2017, preprint, ([arXiv:1708.01531](https://arxiv.org/abs/1708.01531))
 925 Elmegreen B. G., Elmegreen D. M., Chromey F. R., Hasselbacher D. A.,
 926 Bissell B. A., 1996, *AJ*, **111**, 2233
 927 Fall S. M., Efstathiou G., 1980, *MNRAS*, **193**, 189
 928 Fisher D. B., Drory N., 2008, *AJ*, **136**, 773
 929 Flaughner B., 2005, *International Journal of Modern Physics A*, **20**, 3121
 930 Freeman P. E., Izbicki R., Lee A. B., Newman J. A., Conselice C. J., Koeke-
 931 moer A. M., Lotz J. M., Mozena M., 2013, *MNRAS*, **434**, 282
 932 Genzel R., et al., 2008, *ApJ*, **687**, 59

- 933 Graham A. W., Driver S. P., 2005, *Publ. Astron. Soc. Australia*, **22**, 118
- 934 Gray M. E., et al., 2009, *MNRAS*, **393**, 1275
- 935 Häußler B., et al., 2007, *ApJS*, **172**, 615
- 936 Häußler B., Barden M., Bamford S. P., Rojas A., 2011, in Evans I. N.,
937 Accomazzi A., Mink D. J., Rots A. H., eds, *Astronomical Society of*
938 *the Pacific Conference Series Vol. 442, Astronomical Data Analysis*
939 *Software and Systems XX*. p. 155
- 940 High F. W., Rest A., Stubbs C., Stalder B., Challis P., 2009, in *American*
941 *Astronomical Society Meeting Abstracts #214*. p. 680
- 942 Huertas-Company M., Rouan D., Tasca L., Soucaïl G., Le Fèvre O., 2008, *1010*
943 *A&A*, **478**, 971
- 944 Huertas-Company M., et al., 2015, *ApJS*, **221**, 8
- 945 Jooee S., 2009, in Jooee S., Marinova I., Hao L., Blanc G. A., eds, *As-*
946 *tronomical Society of the Pacific Conference Series Vol. 419, Galaxy*
947 *Evolution: Emerging Insights and Future Challenges*. p. 204
- 948 Jooee S., et al., 2009, *ApJ*, **697**, 1971
- 949 Kauffmann G., White S. D. M., Heckman T. M., Ménard B., Brinchmann
950 J., Charlot S., Tremonti C., Brinkmann J., 2004, *MNRAS*, **353**, 713
- 951 Kormendy J., 1977, *ApJ*, **217**, 406
- 952 Kormendy J., 1979, *ApJ*, **227**, 714
- 953 Kormendy J., Kennicutt Jr. R. C., 2004, *ARA&A*, **42**, 603
- 954 Kuutma T., Tamm A., Tempel E., 2017, *A&A*, **600**, L6
- 955 Lahav O., 1995, *Astrophysical Letters and Communications*, **31**, 73
- 956 Lahav O., et al., 1995, *Science*, **267**, 859
- 957 Lani C., et al., 2013, *MNRAS*, **435**, 207
- 958 Law D. R., Steidel C. C., Erb D. K., Pettini M., Reddy N. A., Shapley A. E.,
959 Adelberger K. L., Simenc D. J., 2007, *ApJ*, **656**, 1
- 960 Lin L., et al., 2004, *ApJ*, **617**, L9
- 961 Lintott C. J., et al., 2008, *MNRAS*, **389**, 1179
- 962 Lotz J. M., Primack J., Madau P., 2004, *AJ*, **128**, 163
- 963 Lotz J. M., Jonsson P., Cox T. J., Primack J. R., 2008a, *MNRAS*, **391**, 1137
- 964 Lotz J. M., et al., 2008b, *ApJ*, **672**, 177
- 965 Martig M., Bournaud F., Teyssier R., Dekel A., 2009, *ApJ*, **707**, 250
- 966 Melchior P., Viola M., 2012, *MNRAS*, **424**, 2757
- 967 Miyazaki S., et al., 2012, *SPIE*, 8446, 84460Z
- 968 Mohr J. J., et al., 2012, *SPIE*, 8451, 84510D
- 969 Morganson E., et al., 2018, preprint, ([arXiv:1801.03177](https://arxiv.org/abs/1801.03177))
- 970 Naab T., Burkert A., 2003, *ApJ*, **597**, 893
- 971 Peng C. Y., Ho L. C., Impey C. D., Rix H.-W., 2002, *AJ*, **124**, 266
- 972 Peng C. Y., Ho L. C., Impey C. D., Rix H.-W., 2010, *AJ*, **139**, 2097
- 973 Petrosian V., 1976, *ApJ*, **209**, L1
- 974 Pieres A., et al., 2017, *MNRAS*, **468**, 1349
- 975 Postman M., et al., 2005, *ApJ*, **623**, 721
- 976 Scarlata C., et al., 2007a, *ApJS*, **172**, 406
- 977 Scarlata C., et al., 2007b, *ApJS*, **172**, 494
- 978 Schade D., Lilly S. J., Crampton D., Hammer F., Le Fevre O., Tresse L.,
979 1995, *ApJ*, **451**, L1
- 980 Schawinski K., et al., 2007, *ApJS*, **173**, 512
- 981 Schawinski K., et al., 2014, *MNRAS*, **440**, 889
- 982 Sellwood J. A., 2014, *Rev. Mod. Phys.*, **86**, 1
- 983 Sérsic J. L., 1963, *Boletin de la Asociacion Argentina de Astronomia La*
984 *Plata Argentina*, **6**, 41
- 985 Sevilla I., et al., 2011, in *Particles and fields. Proceedings, Meeting of*
986 *the Division of the American Physical Society, DPF 2011, Providence,*
987 *USA, August 9-13, 2011.* ([arXiv:1109.6741](https://arxiv.org/abs/1109.6741)), http://lss.fnal.gov/cgi-bin/find_paper.pl?conf-11-525
- 988 Simard L., Mendel J. T., Patton D. R., Ellison S. L., McConnell A. W.,
989 2011, *ApJS*, **196**, 11
- 990 Simmons B. D., et al., 2017, *MNRAS*, **464**, 4420
- 991 Soumagnac M. T., et al., 2015, *MNRAS*, **450**, 666
- 992 The DES Collaboration 2016, *MNRAS*, **460**, 1270
- 993 Toomre A., 1977, in Tinsley B. M., Larson D. Campbell R. B. G., eds,
994 *Evolution of Galaxies and Stellar Populations*. p. 401
- 995 Toomre A., Toomre J., 1972, *ApJ*, **178**, 623
- 996 Weinmann S. M., van den Bosch F. C., Yang X., Mo H. J., 2006, *MNRAS*,
997 **366**, 2
- 998 White S. D. M., Rees M. J., 1978, *MNRAS*, **183**, 341
- 1000 Willett K. W., et al., 2017, *MNRAS*, **464**, 4176
- 1001 Woo J., Dekel A., Faber S. M., Koo D. C., 2015, *MNRAS*, **448**, 237
- 1002 York D. G., et al., 2000, *The Astronomical Journal*, **120**, 1579
- 1003 Zamojski M. A., et al., 2007, *ApJS*, **172**, 468
- 1004 de Jong R. S., 1996, *A&A*, **313**, 45
- 1005 van der Wel A., 2008a, in Bureau M., Athanassoula E., Barbuy B., eds, *IAU*
1006 *Symposium Vol. 245, Formation and Evolution of Galaxy Bulges*. pp
1007 59–62, [doi:10.1017/S1743921308017286](https://doi.org/10.1017/S1743921308017286)
- 1008 van der Wel A., 2008b, *ApJ*, **675**, L13
- 1009 van der Wel A., et al., 2012, *The Astrophysical Journal Supplement Series*,
1010 **203**, 24

APPENDIX A: CALIBRATION MAPS FOR THE *G* AND *R* FILTERS

In this Appendix we present the calibration maps for both parametric and non-parametric measurements in the *g* and *r* bands. They were obtained following the procedure described in Sections 2.2 and 5.4 for parametric and non-parametric fits, respectively. The maps are displayed following the same conventions adopted for visualising the calibration maps in the *i* band. Those maps are shown in Figures 9 and 13.

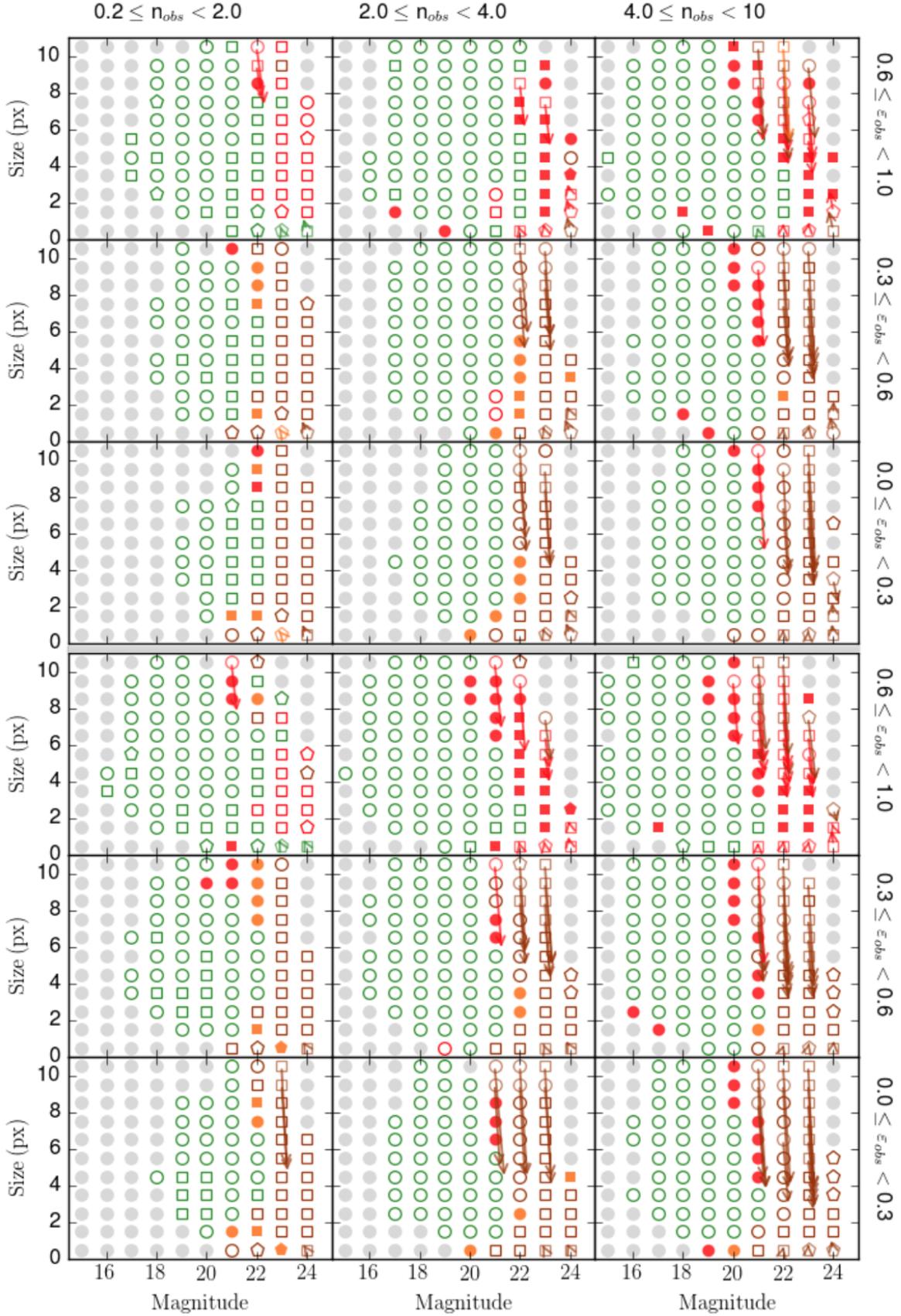


Figure A1. Map of the corrections for Sérsic parameters in the g (upper panel) and r (lower panel) filters, obtained through the simulation routine described in Section 2.2. Symbols and colours have the same meaning as Figure 9.

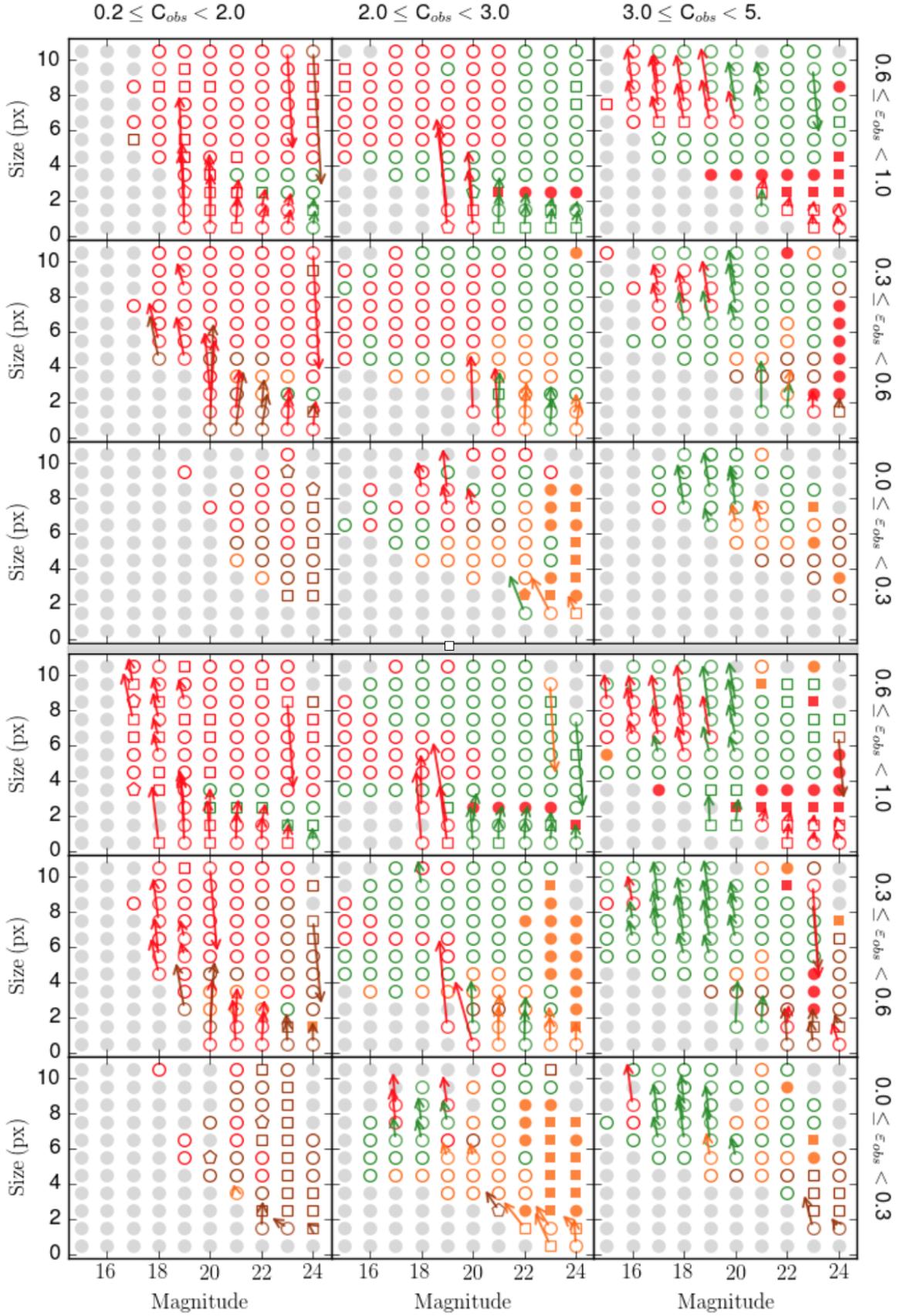


Figure A2. Map of the corrections for ZEST+ output in the g (upper panel) and r (lower panel) filters, obtained through the simulation routine described in Section 5.4. Symbols and colours have the same meaning as Figure 13.

1020 **APPENDIX B: CATALOG MANUAL**

1021 A description of the columns of the catalogue follows, both for parametric and non-parametric fits. In order to distinguish between filters, the
1022 parameters can be labelled with $_X$, where $X = g, r, i$.

1023 **B1 Identification columns**

1024 COADD_OBJECT_ID - Identifier assigned to each object in the co-add DES Y1 dataset, reported here from the *Gold Catalogue*.

1025 TILENAME - Column reporting the name of the tile image where the object lies.

1026 ID - Rows enumerator, running for 1 to the total number of entries in the catalogue.

1027 RA - Right Ascension from the Y1A1 GOLD catalogue.

1028 DEC - Declination from the Y1A1 GOLD catalogue.

1029

1030 **B2 SExtractor parameters for star-galaxy separation and signal-to-noise**

1031 SG - Linear combination of the star-galaxy classifier SPREAD_MODEL and its uncertainty, SPREADERR_MODEL, according to Equation 4.
1032 A cut in $SG > 0.005$ is recommended.

1033 SN_X - Signal-to-noise expressed as the ratio between FLUX_AUTO_X and FLUXERR_AUTO_X.

1034 **B3 Columns for Parametric Fits**

1035 *B3.1 Selection and pre-fitting classification flags*

1036 SELECTION_FLAGS_X - If equal to 1, then the relative object has been selected, according to the requirements described in Section 3.1.
1037 It can assume other numerical values in the following cases:

- 1038 • if the object passes the selection requirements, but is not included in the intersection between the DESDM catalogues and the Y1A1
- 1039 GOLD catalogue, then this flag is set to 2;
- 1040 • if the object passes the selection requirements, but it is fainter than $GOLD_MAG_AUTO_i = 23$, then the flag is set to 3;
- 1041 • if the object enters in the previous category, but it has no match with the Y1A1 GOLD catalogue, then the flag is set to 4.

1042 If the object is not selected because it doesn't pass any of the selection requirements, then the SELECTION_FLAGS_X and all the other flags
1043 are set to zero.

1044 The catalogue version made available to the users includes all the objects which have been selected at least in one of the three bands g, r, i .

1045

1046 C_FLAGS_X - Number of neighbours in the fitted stamp.

1047 MAX_OVERLAP_PERC_X - Percentage of the central galaxy isophotes overlapping with the closest neighbour. If there are no neighbours
1048 or no overlapping neighbours, then it is set to 0. A cut in $MAX_OVERLAP_PERC_X < 50$ is recommended.

1049 *B3.2 Parametric measurements (GALFIT)*

MAG_SERSIC_X - GALFIT value for the magnitude of the galaxy. The value already includes the calibration listed in the column
MAG_CAL_X.

RE_SERSIC_X - GALFIT measure of the half light radius (or Effective radius) of the galaxy. It is expressed in pixels and is already
calibrated. The correction is reported in the column RE_CAL_X.

N_SERSIC_X - GALFIT output for the Sérsic Index. The measure is calibrated, and the can find the relative correction in the column
N_SERSIC_CAL_X.

ELLIPTICITY_SERSIC_X - Ellipticity of the galaxy, calculated by subtracting from unity the GALFIT estimate for the axis-ratio. The
value is corrected and the calibration is accessible through the column ELLIPTICITY_SERSIC_CAL_X.

OUTLIERS_X - If equal to 1, it labels the objects classified as outliers in the catalogue validation process.

FIT_STATUS_X - If equal to 1, this flag selects all the objects with a successfully validated and calibrated converged fit.

Important note: by applying the recommended cut $FIT_STATUS_X = 1$, the user is able to collect the sample of validated and calibrated
objects in the X filter. This cut is equivalent to applying all together the cuts which are recommended in terms of sample selection, fitting
convergence, bad regions masking, exclusion of outliers and significantly overlapping objects, minimization of stellar contamination. A

summarising scheme follows:

$$(\text{FIT_STATUS_X}=1) = \left\{ \begin{array}{l} \text{FLAGS_BADREGION}=0 \\ \text{SG}>0.005 \\ \text{SELECTION_FLAGS_X}=1 \\ \text{FIT_AVAILABLE_X}=1 \wedge \text{WARNING_FLAGS_CENTRAL_X}=0 \\ \text{MAX_OVERLAP_PERC_X}<50 \\ \text{OUTLIERS_X}=0 \\ \text{PARAMETER_CAL_X}<99, \end{array} \right.$$

1050 where the voice `PARAMETER_CAL_X` can be `MAG_CAL_X` etc. In absence of calibration the correction value is set to 99.
 1051 For a cleaner sample the user can associate the cut in `FIT_STATUS_X` to the condition `SN_X>30`.

1052 **B4 Columns for non-parametric coefficients (ZEST+)**

- 1053 `SELECTION_NP_X` - If equal to 1, the object is selected in the X filter, otherwise it is 0.
 1054 `FIT_STATUS_NP_X` - If equal to 1, this flag selects all the objects with successfully validated and calibrated measurements.
 1055 `CONCENTRATION_X` - ZEST+ measurement for the Concentration of light. See Equation 11 for its definition.
 1056 `ASYMMETRY_X` - ZEST+ value for the Asymmetry (see Equation 15).
 1057 `CLUMPINESS_X` - ZEST+ value for the Clumpiness (see Equation 16).
 1058 `GINI_X` - Measure of the Gini parameter, defined in Equation 17.
 1059 `M20_X` - Measure of the M20 parameter, for more details see Equation 19.

1060

1061 This paper has been typeset from a $\text{T}_{\text{E}}\text{X}/\text{L}^{\text{A}}\text{T}_{\text{E}}\text{X}$ file prepared by the author.